

X-ray Line Profile Diagnostics of Shock Heated Stellar Winds

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Hot Star Winds

Chandra HETGS

- Nested parabolic mirrors
- High Energy Transmission Grating:
 - Effective area $\sim 10 \text{ cm}^2$
 - Resolution ~ 1000

Hot stars have massive, highly supersonic radiation driven winds:

- Observed in UV absorption lines
- Velocities of order few 1000 km s^{-1}
- Densities of order 10^{10} cm^{-3}
- Mass-loss rates up to $10^{-5} M_{\text{sun}} \text{ yr}^{-1}$
- Steady-state models based on radiation pressure are quantitatively successful
- But many indications of time-variability in hot star winds: Shock heating and possibly some connection to photospheric variability and magnetic fields

The heating mechanisms, not to mention the physical properties, of the X-ray emitting plasma on these hot stars is not known. Leading theories:

1. Line force instability generated shocks, leading to hot plasma distributed throughout the wind (described by a filling factor),
2. Magnetically confined wind shocks,
3. Solar-type coronal magnetic heating (but hot stars are not thought to have dynamos and coronae).

In all cases, the X-ray emitting plasma is thermal and optically thin, emitting photons in lines. Especially in case (1) there is a bulk, cold wind component (which leads to the UV absorption lines) that is a source of X-ray continuum opacity.

- The speed of these winds has the potential to produce substantial Doppler broadening of the lines, while the continuum absorption can alter the line shapes through the spatial dependence of the absorption coupled with the spatial dependence of the Doppler shift in an organized flow.

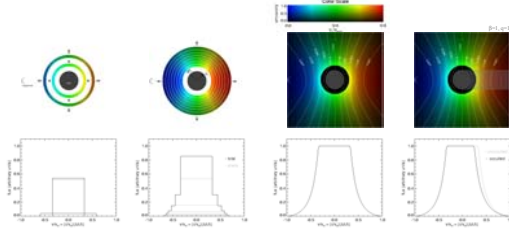
- High-resolution X-ray spectroscopy of astrophysical objects has only in the last two years become feasible, and hot stars are one of the very few types of sources for which telescopes like Chandra can resolve X-ray line shapes.

- We can use the observed line shapes to infer the spatial- and velocity-distributions of the X-ray emitting plasma on hot stars, and thereby constrain models of X-ray production and wind dynamics.

Line Transport in an Expanding and Continuum-Absorbing Medium

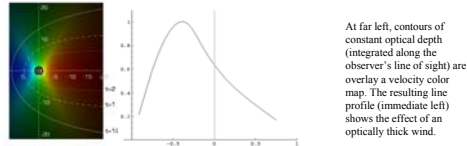
Spherically Symmetric Models with Absorption

Doppler shifting of the expanding wind broadens lines. To the observer on the left, the front of the wind is blueshifted and the back is redshifted.



The wind is depicted spatially in the color plots, with the hue indicating velocity with respect to an observer on the left, and the brightness of the ink indicating emissivity (scaling as density squared). Note the color scale above the third panel. Under each image is the resulting line profile with the bluest wavelengths (expressed in velocity units) on the left, and the reddest on the right.

Including continuum absorption by the cold component of the wind also preferentially removes red photons.



At far left, contours of constant optical depth (integrated along the observer's line of sight) are overlay a velocity color map. The resulting line profile (immediate left) shows the effect of an optically thick wind.

Phenomenological model of wind emission and absorption with four parameters.

We have developed a physically meaningful line-profile model, yet one that is simple and not tied to any one proposed mechanism of hot-star X-ray production. Described in Owocki & Cohen (2001, *ApJ*, 559, 1108), the model assumes a smoothly and spherically symmetrically distributed accelerating X-ray emitting plasma subject to continuum attenuation by the cold stellar wind.

$$L_x = 8\pi^2 \int_0^1 \int_0^\pi \int_0^{2\pi} d\mu \int_R^{\infty} r^2 \eta_c(\mu, r) e^{-\tau(r, \lambda)} dr$$

All the emission physics is hidden in the emissivity, η . Note that spherical coordinates (μ, r) are natural for the symmetry of the wind emission.

The velocity is assumed to be of the form $v(r) = v_\infty(1 - R_0/r)^\beta$

With the observer looking at the star and wind from one side, cylindrical coordinates (ρ, z) are more natural.

$$\tau(\rho, z) = \tau_* \int_0^\infty \frac{R_0 dz'}{r'^2 (1 - R_0/r')^\beta} \quad \text{where} \quad \tau_* \equiv \frac{\kappa \dot{M}}{4\pi v_\infty R_*}$$

τ_* and β parameterize the absorption. And $r' = \sqrt{\rho^2 + z'^2}$

$$\eta(r) \sim \delta(\lambda - \lambda_\alpha (1 - \mu v/c)) f(r) \sim \delta(\lambda - \lambda_\alpha) e^{-q} \quad \text{for } r \geq R_0 \quad q \text{ and } R_0 \text{ parameterize the radial X-ray filling factor, thus the emissivity.}$$

It is this **delta function** that allows us to map μ, r into wavelength, λ . We solve these equations numerically with *Mathematica*.

The Line Profile Model is General

It can parameterize many different types of wind X-ray distributions, allowing for the testing of different physical theories

