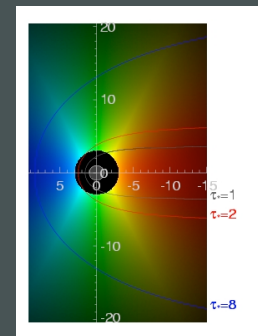
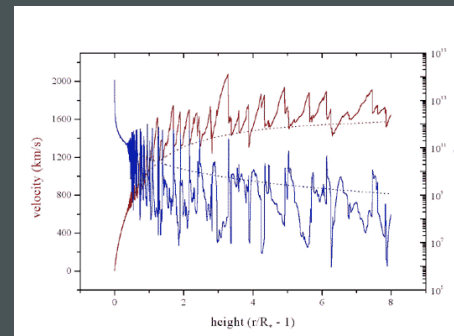
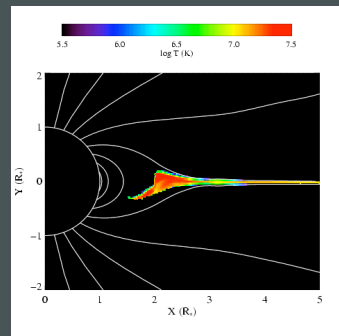
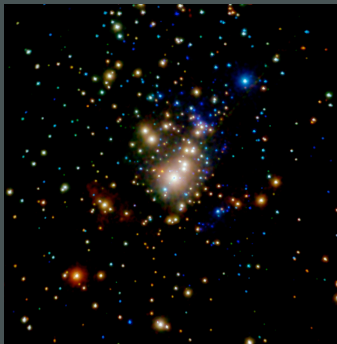


# X-ray Spectroscopy of the Radiation-Driven Winds of Massive Stars: Line Profile and Line Ratio Diagnostics

David Cohen  
Swarthmore College





Carina/Keyhole Nebula (HST)

massive stars:

20 to 100  $M_{\text{sun}}$

$10^6 L_{\text{sun}}$

$T \sim 50,000 \text{ K}$

# Chandra X-ray Telescope image of the Orion Nebula Cluster

young, massive star:  
 $\theta^1$  Ori C



Color coded according to photon energy (red: <1keV; green 1 to 2 keV; blue > 2 keV)

## Keyhole Nebula



Hubble  
Heritage

NASA and The Hubble Heritage Team (STScI) • Hubble Space Telescope WFPC2 • STScI-PRC00-06



Whirlpool/M51 (HST)

# 1000 yr old supernova remnant



Crab Nebula (WIYN)

wind-blown bubble: stellar wind impact on its environment



NGC 6888 Crescent Nebula (Tony Hallas)

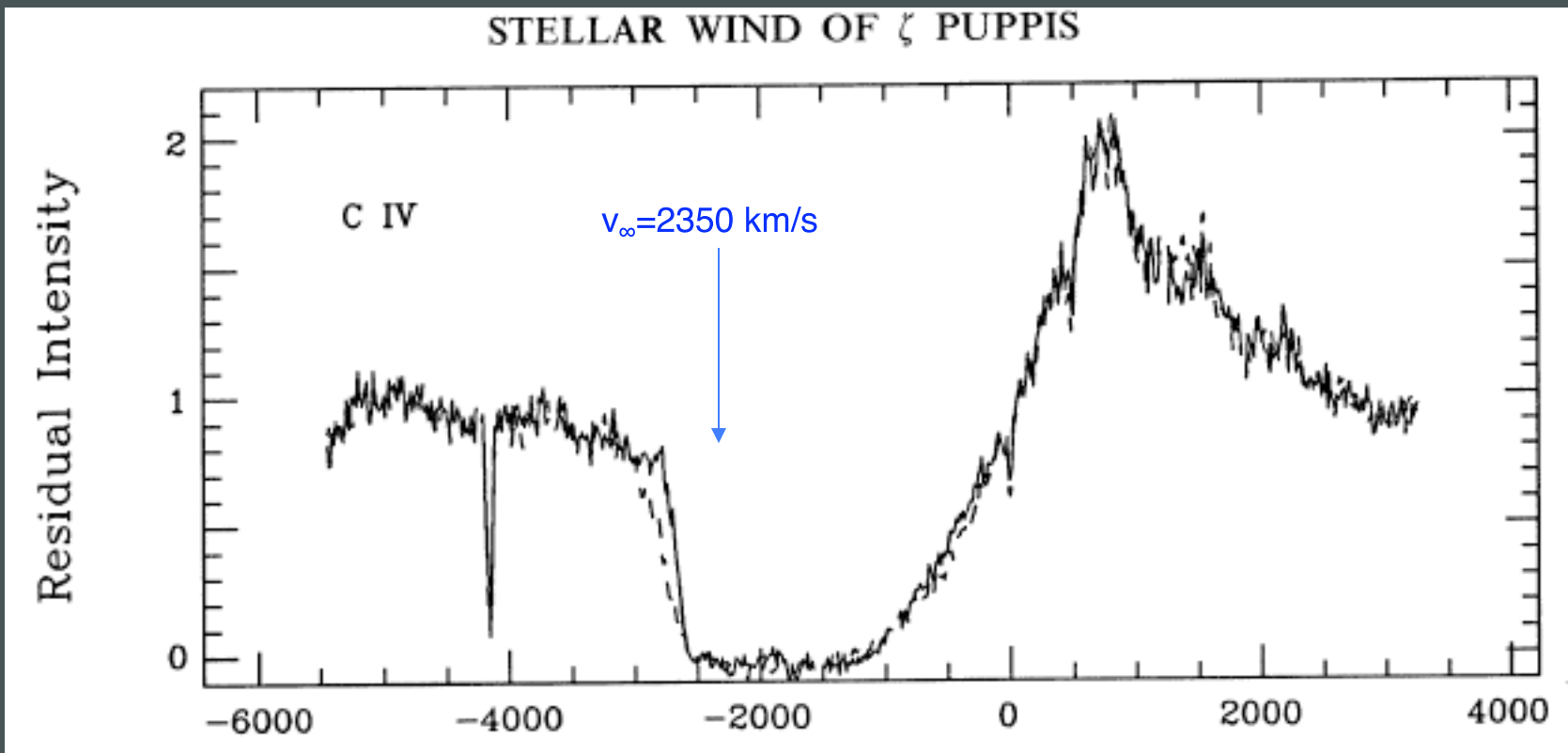
# Radiation-driven massive star winds

$$\dot{M} \sim 10^{-6} M_{\text{sun}}/\text{yr}$$



UV spectrum: C IV 1548, 1551 Å

STELLAR WIND OF ζ PUPPIS



Velocity (km/s)



Power in these winds:

$$\frac{1}{2} \dot{M} v_{\infty}^2 \approx 3 \times 10^{36} \text{ erg s}^{-1}$$
$$\approx .001 L_*$$

$$L_{\text{sun}} = 4 \times 10^{33} \text{ erg s}^{-1}$$

$$L_{\text{massive}} \approx 4 \times 10^{39}$$

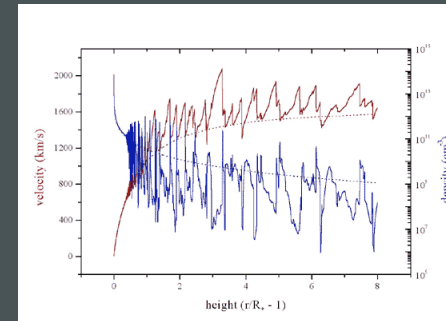
while the x-ray luminosity

$$L_X \approx 10^{-7} L_*$$

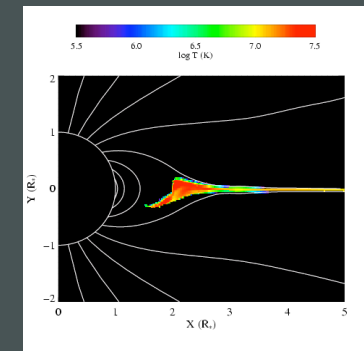
To account for the x-rays, only **one part in  $10^{-4}$**  of the wind's mechanical power is needed to heat the wind

# Three models for massive star x-ray emission

1. Instability driven shocks



2. Magnetically channeled wind shocks

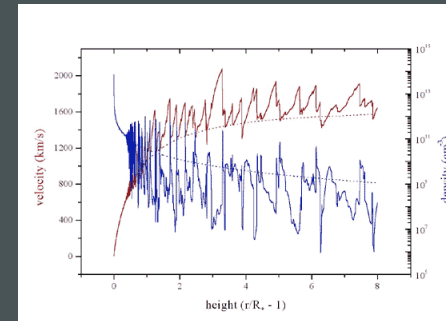


3. Wind-wind interaction in close binaries

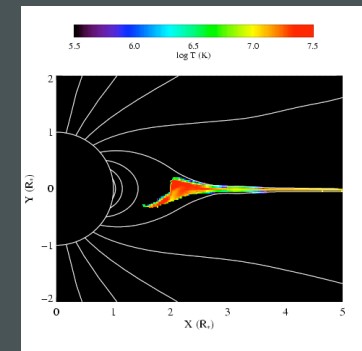


# Three models for massive star x-ray emission

1. Instability driven shocks



2. Magnetically channeled wind shocks



3. Wind-wind interaction in close binaries

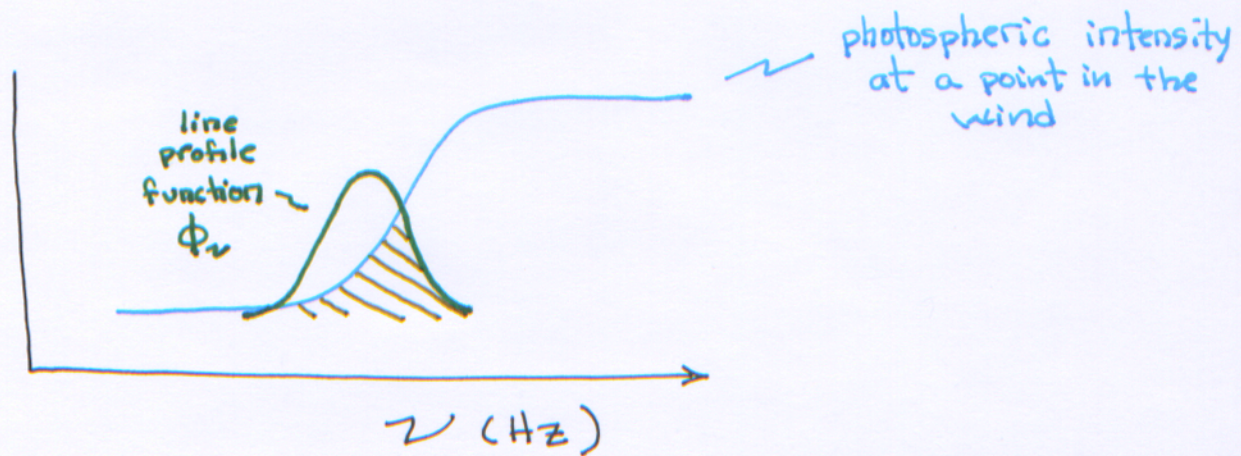




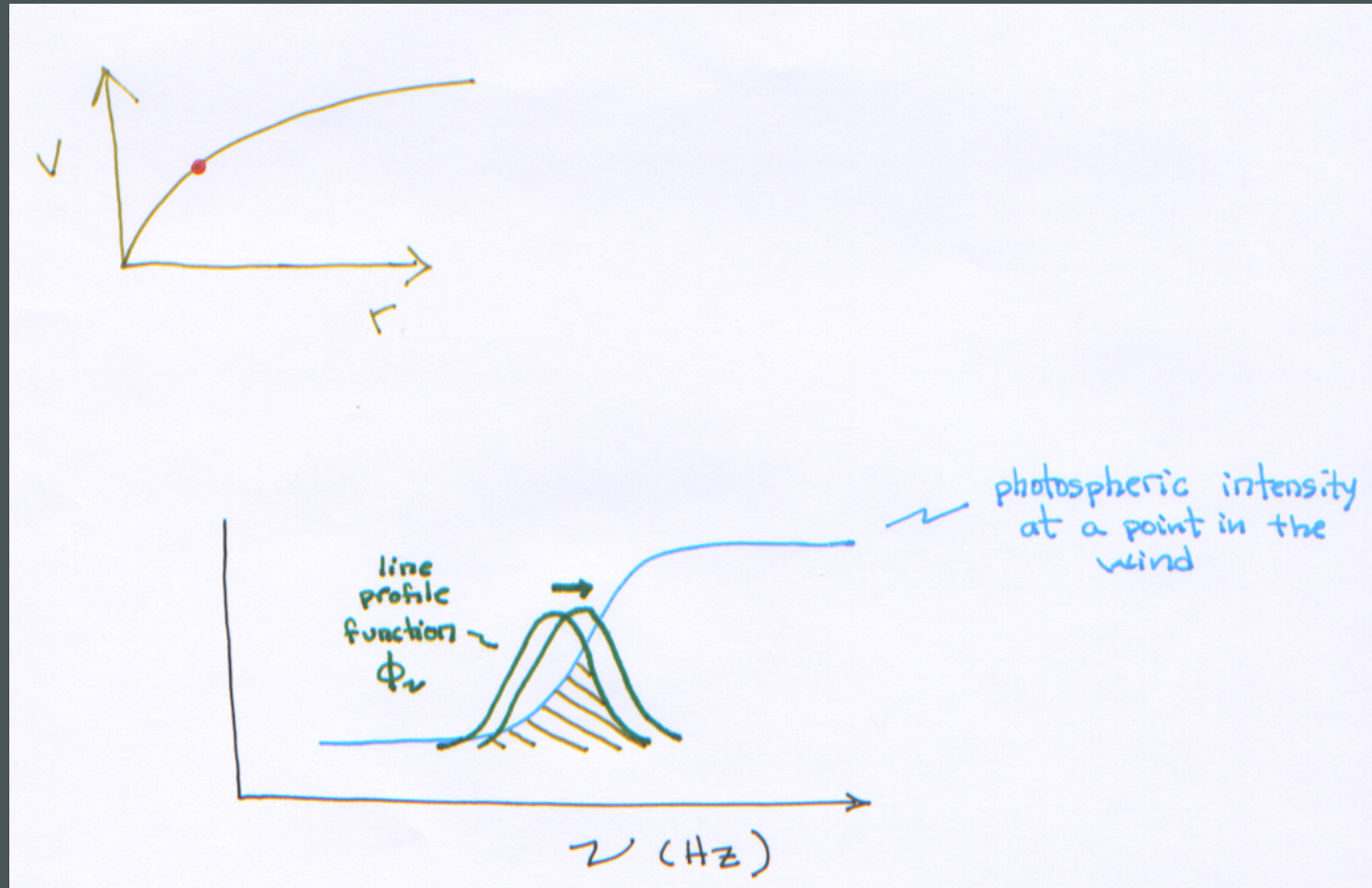
# The Line-Driven Instability (LDI; Milne 1926)



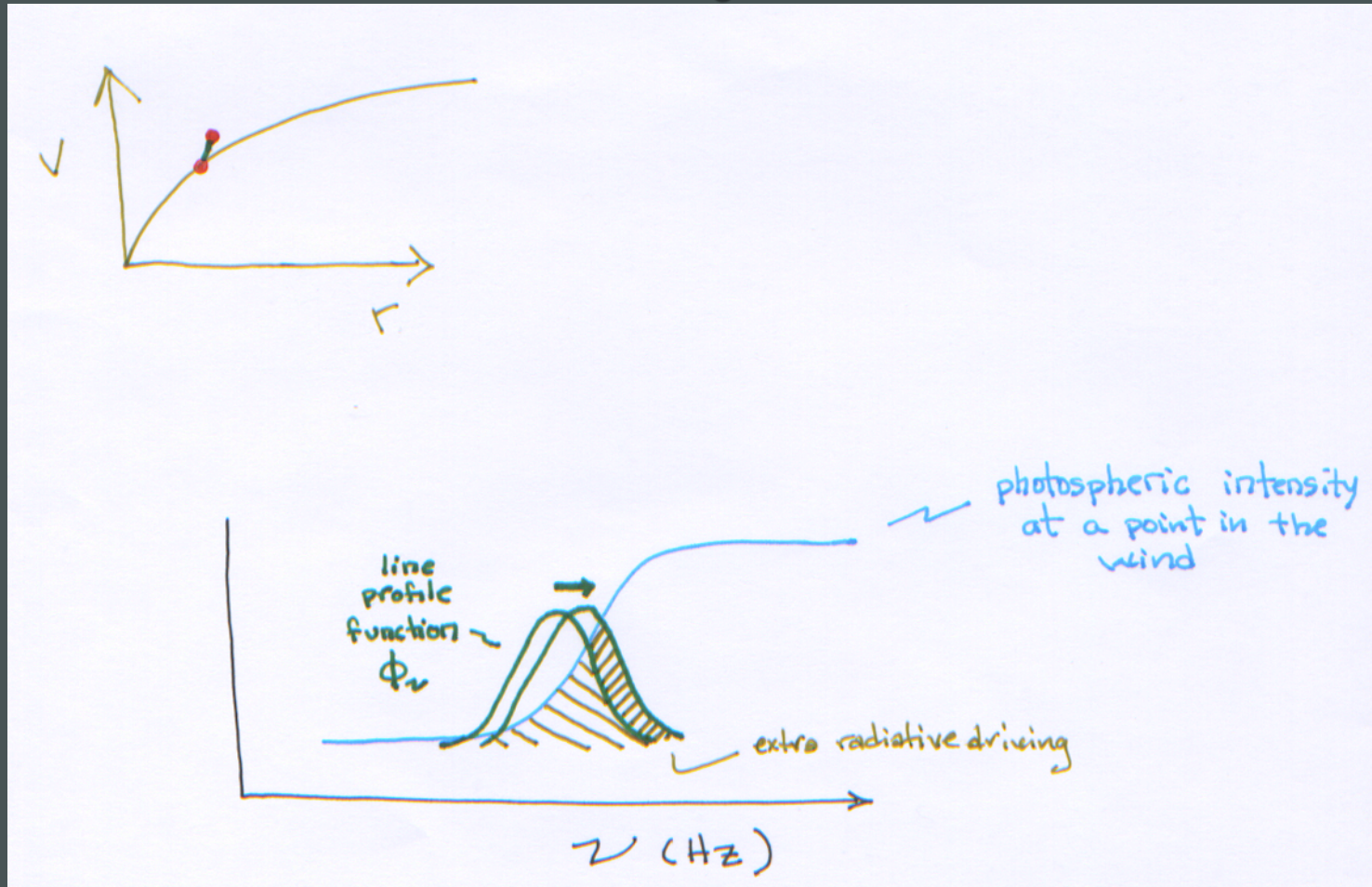
$$v = v_{\infty} (1 - R_*/r)^{\beta}$$



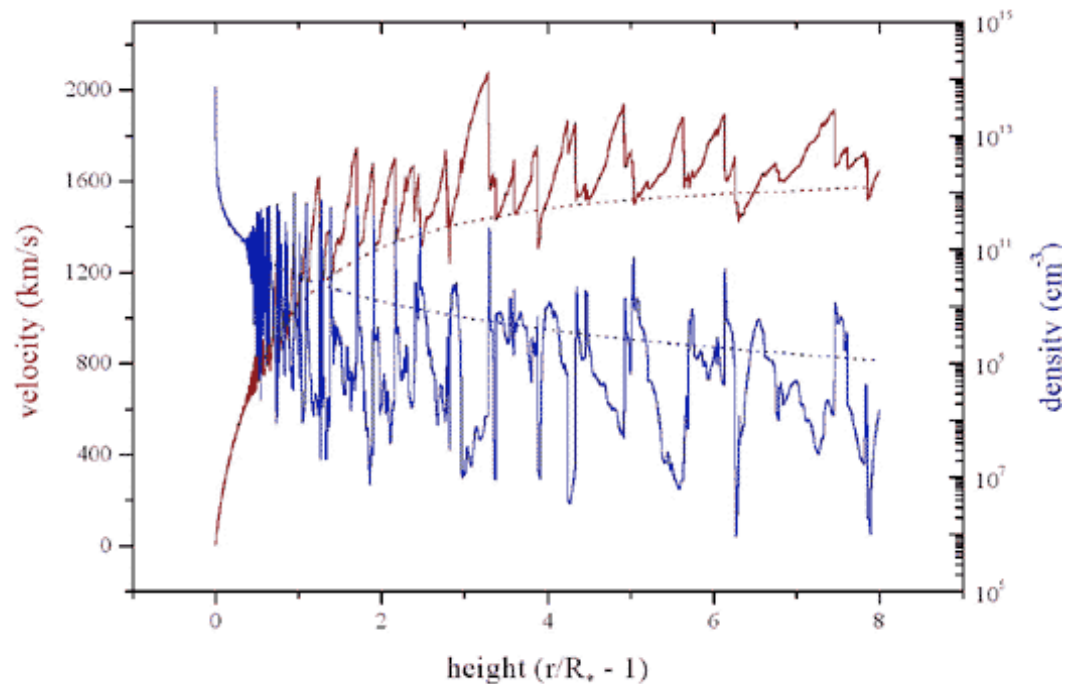
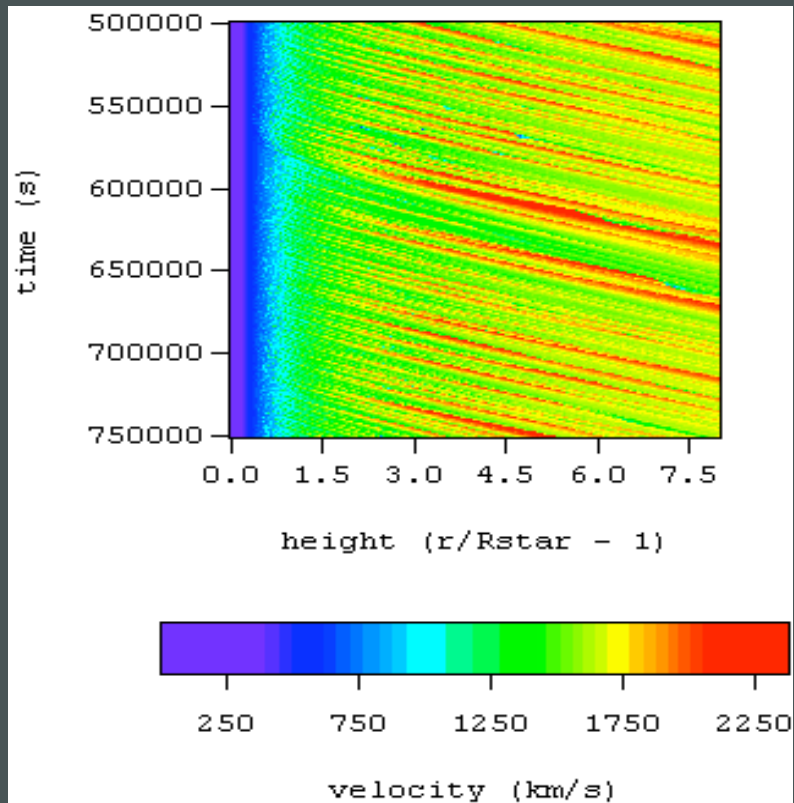
Consider a positive velocity perturbation



Positive feedback: ion moves out of the Doppler shadow, sees more radiation, gets accelerated...



# 1-D rad-hydro simulation of a massive star wind

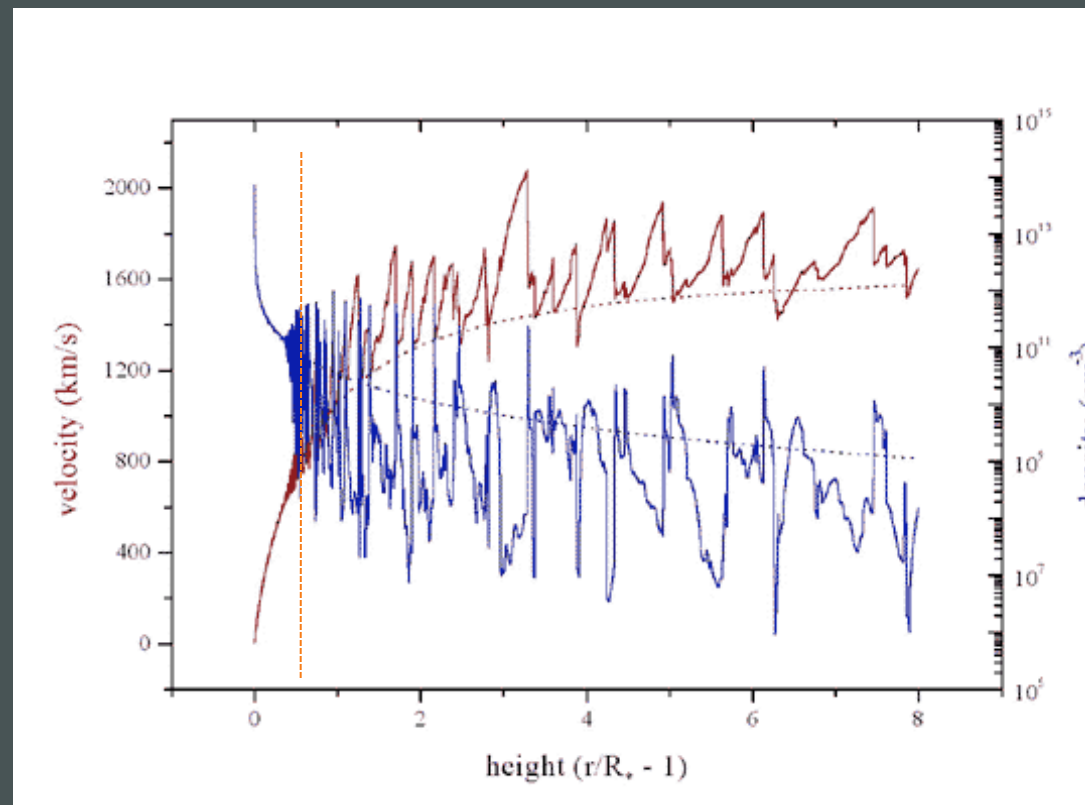


Radiation line driving is inherently unstable:  
shock-heating and X-ray emission  
Owocki, Castor, & Rybicki 1988

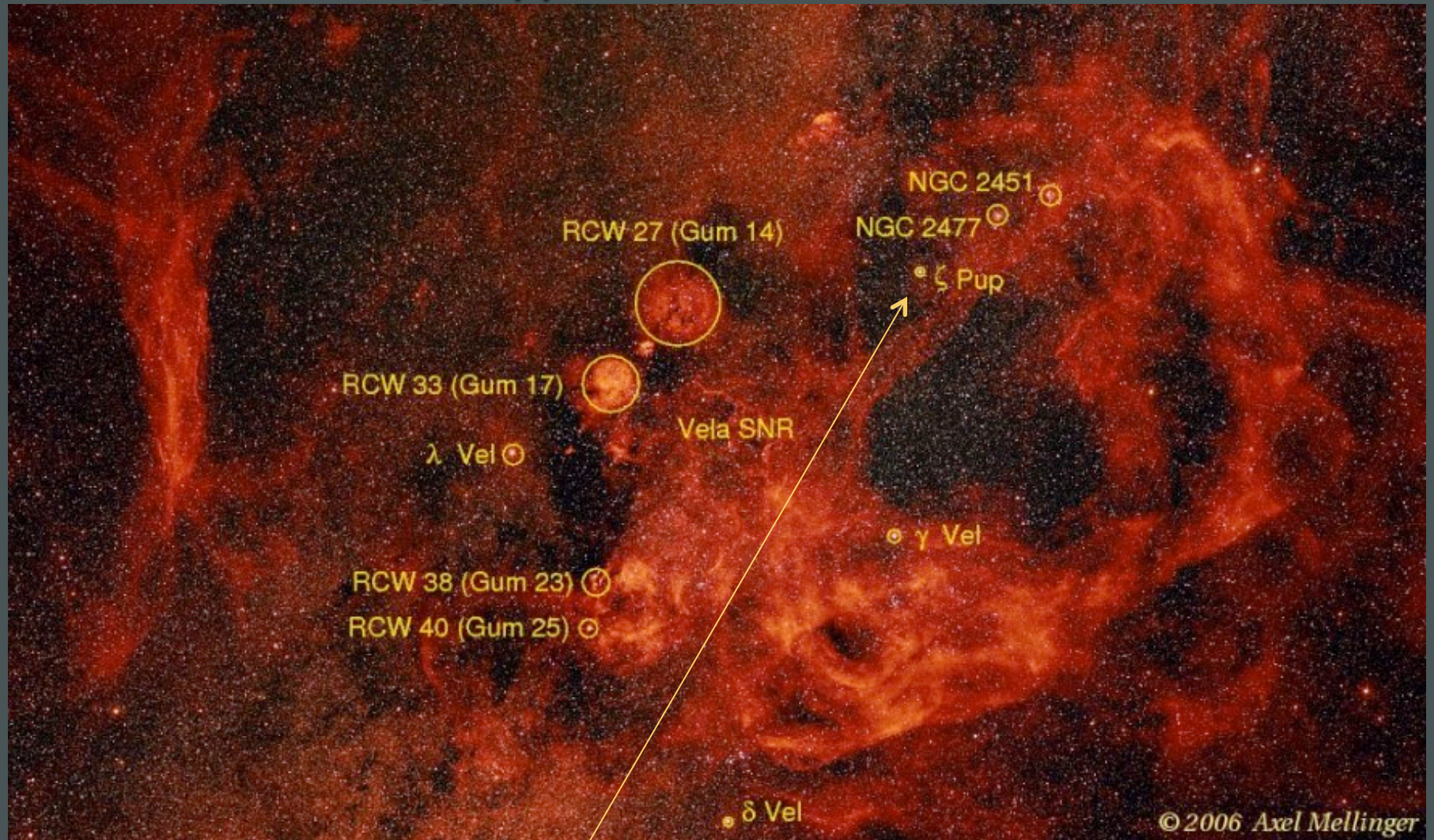


# Predictions of the rad-hydro wind simulations:

1. Significant Doppler broadening of x-ray emission lines due to bulk motion of the wind flow (1a. Shock onset several tenths  $R_*$  above the surface)
2. Bulk of the wind is cold and unshocked – source of attenuation of the X-rays.



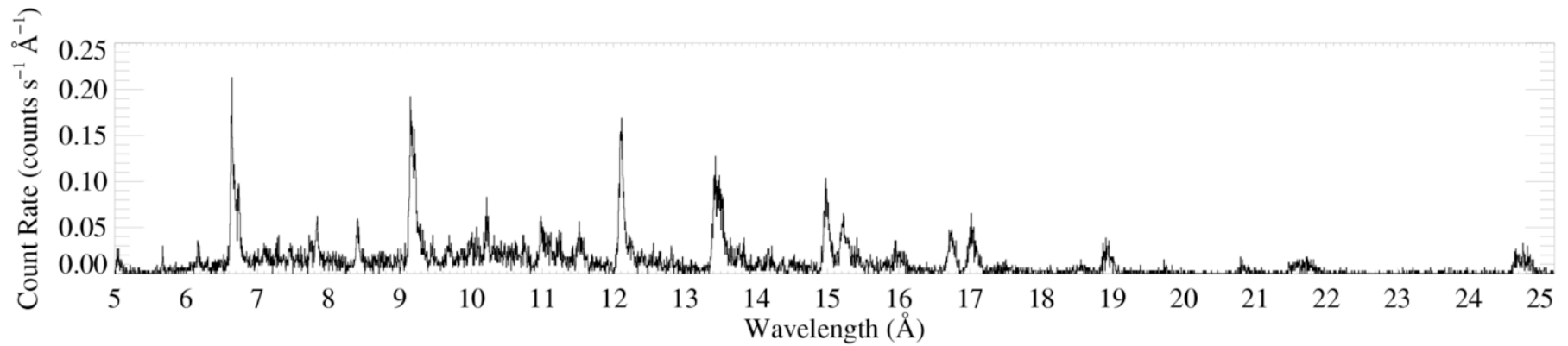
## $\zeta$ Puppis in the Gum Nebula



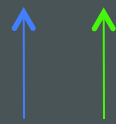


*Chandra HETGS/MEG spectrum*  
( $R \sim 1000 \sim 300 \text{ km s}^{-1}$ )

$\zeta$  Pup



Si



Mg



Ne



Fe

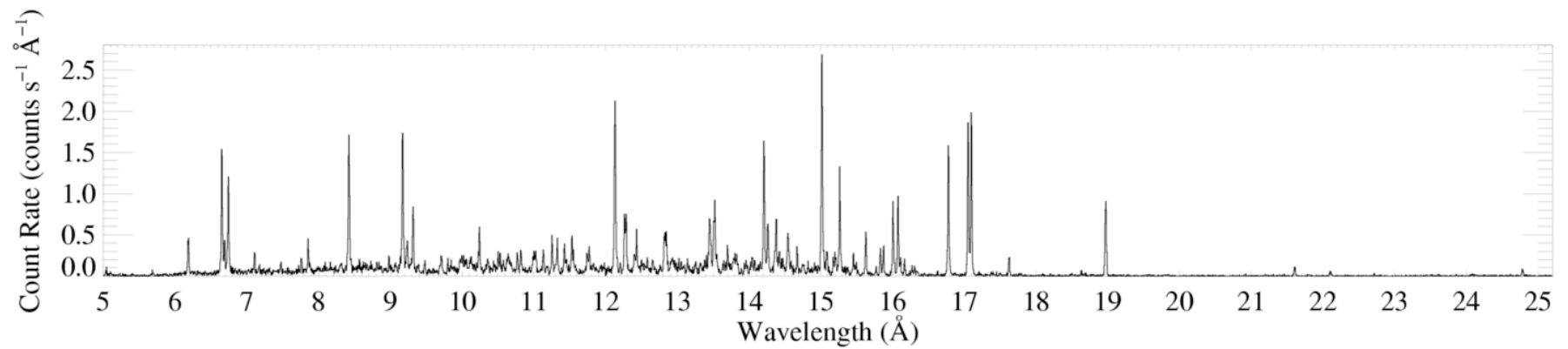
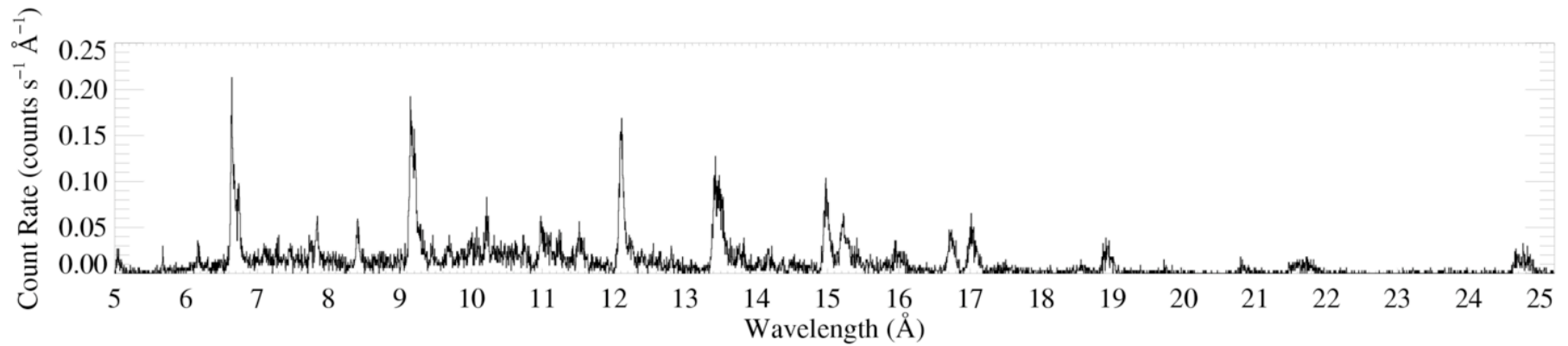


O



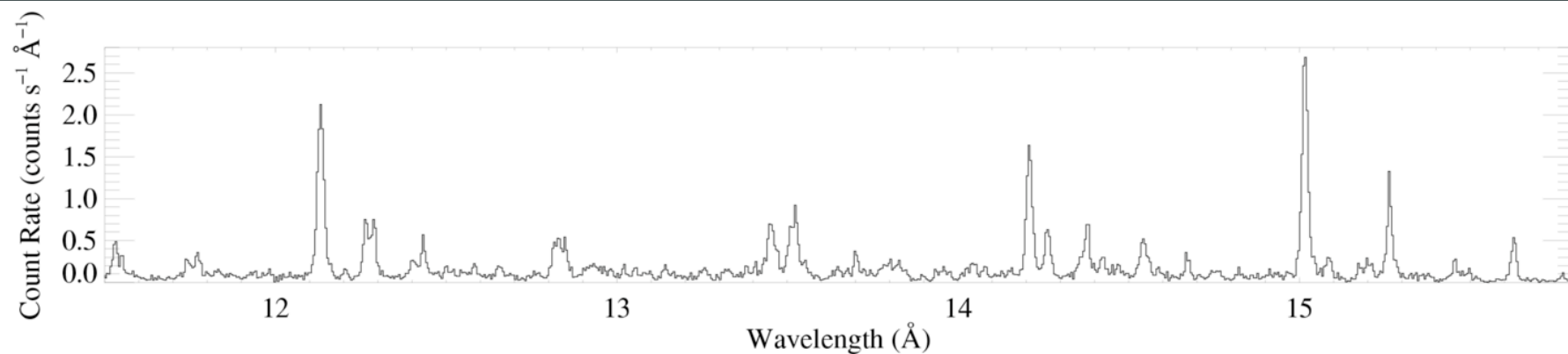
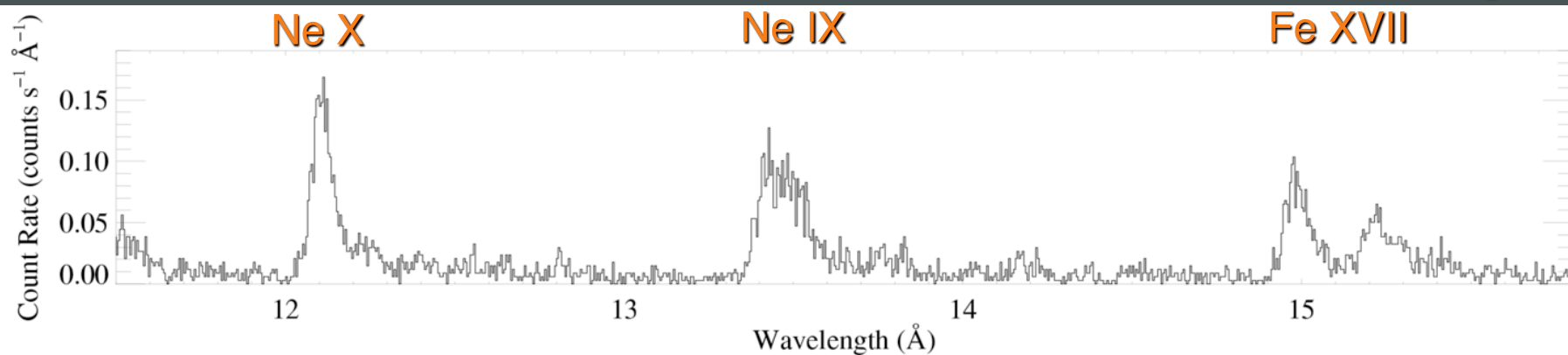
H-like  
He-like

$\zeta$  Pup

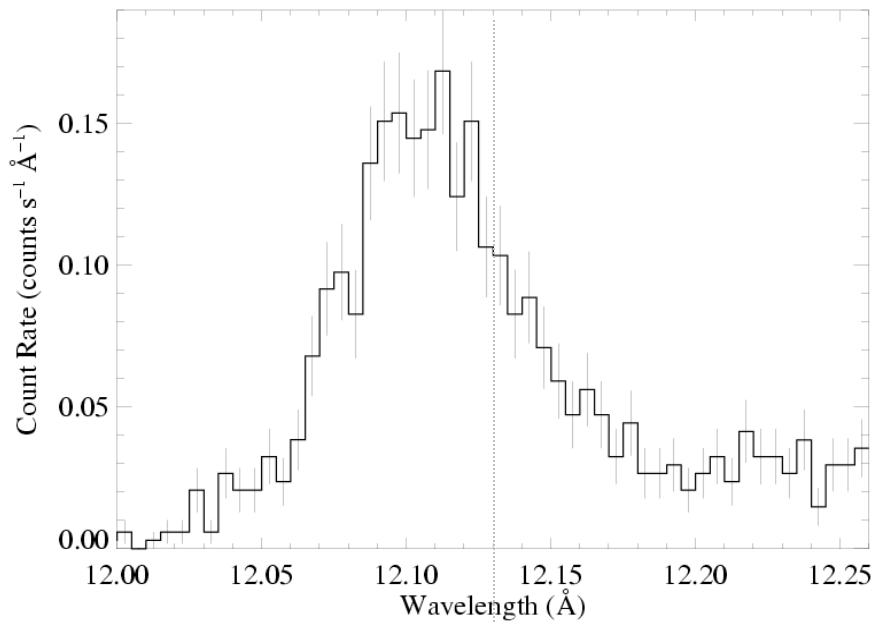


Low-mass star (Capella) for comparison

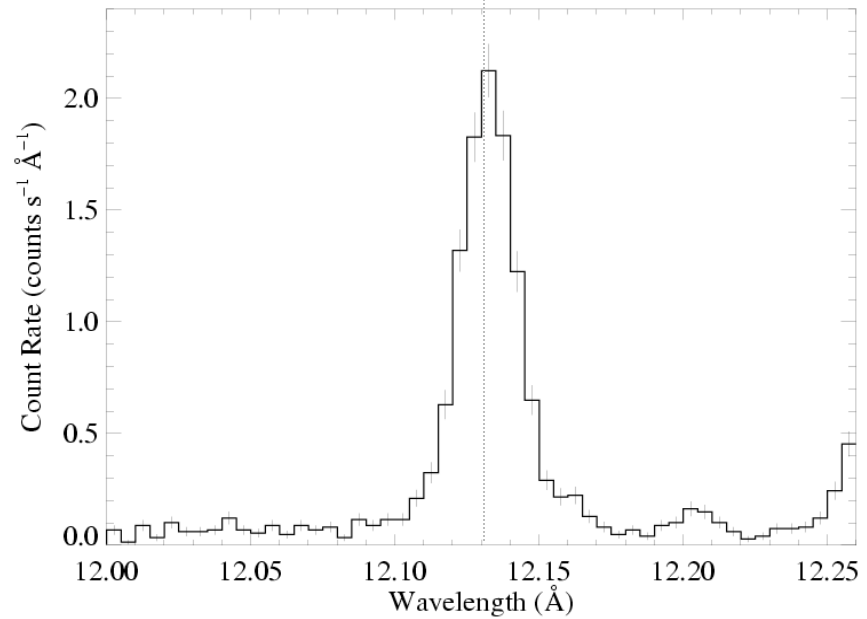
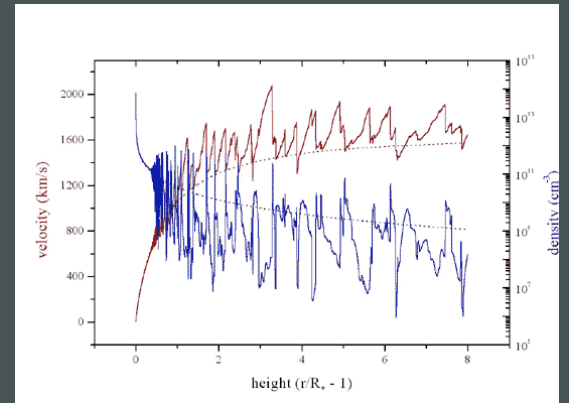
$\zeta$  Pup



Capella



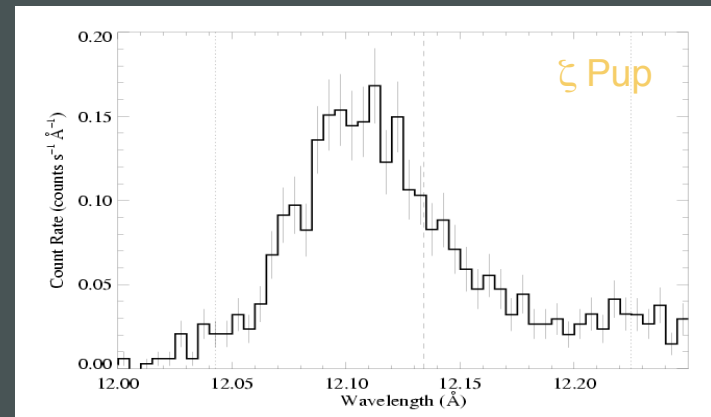
ζ Pup  
massive



Capella  
low mass



The x-ray emission lines are broad: agreement with rad hydro simulations

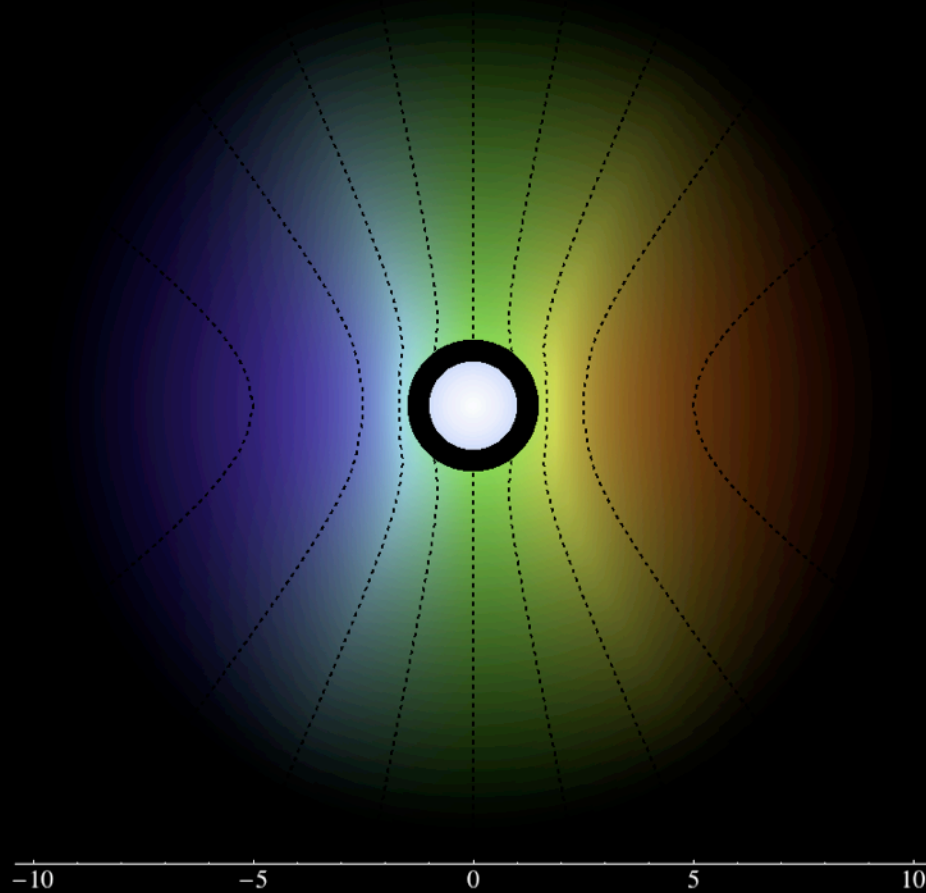


But... they're also blue shifted and asymmetric  
Is this predicted by the wind shock scenario?



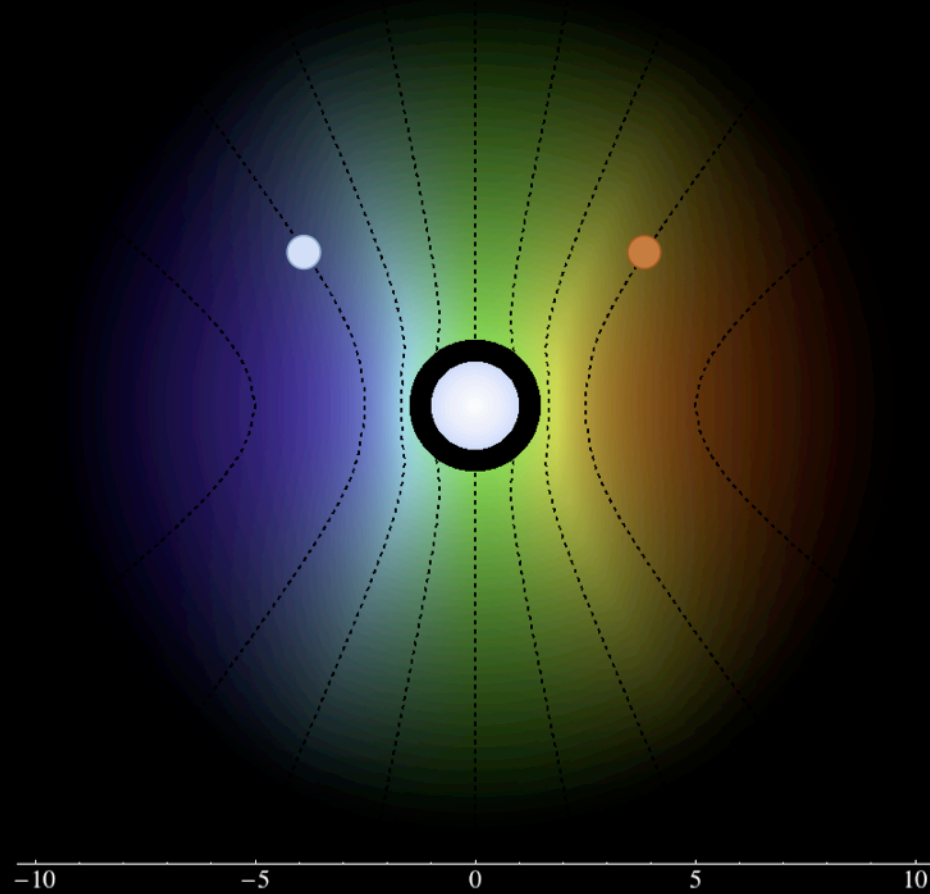
# Wind Profile Model

A

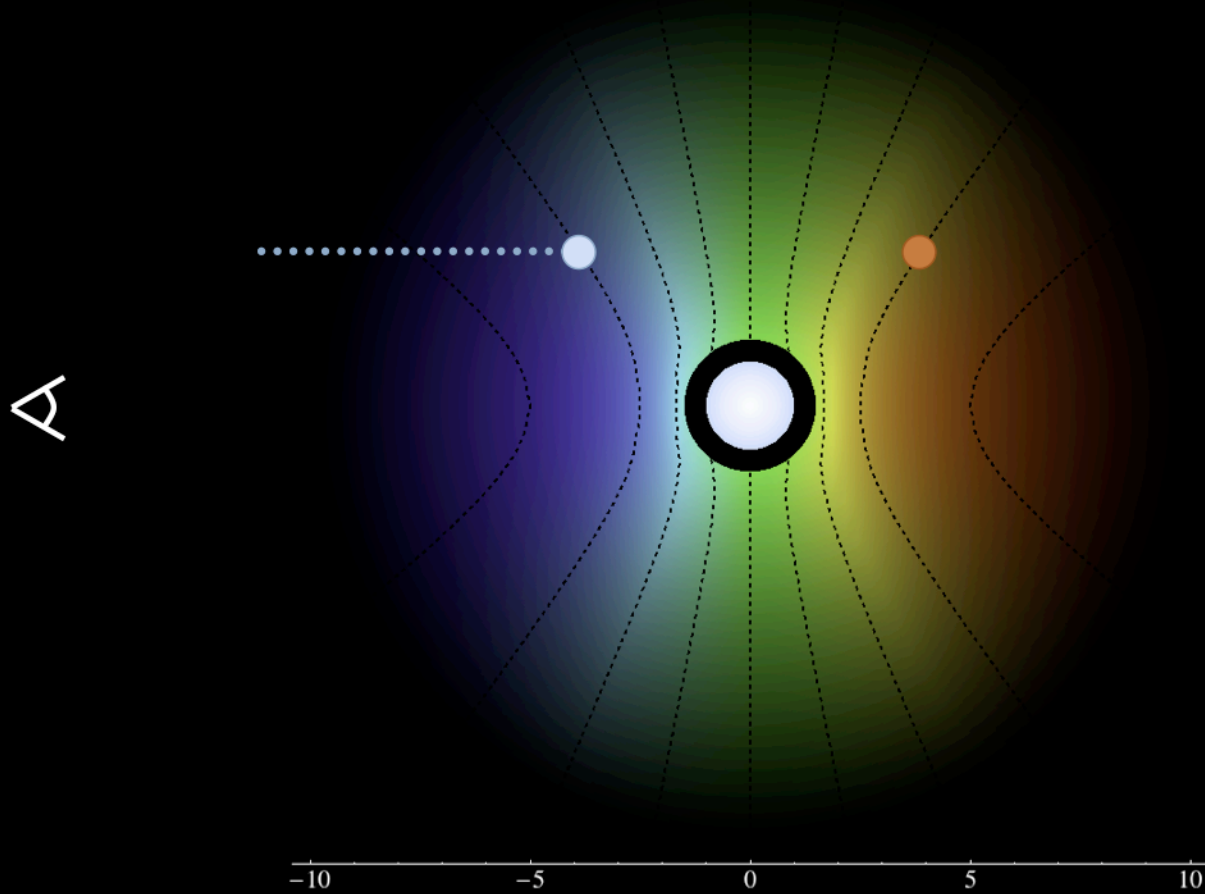


# Wind Profile Model

A

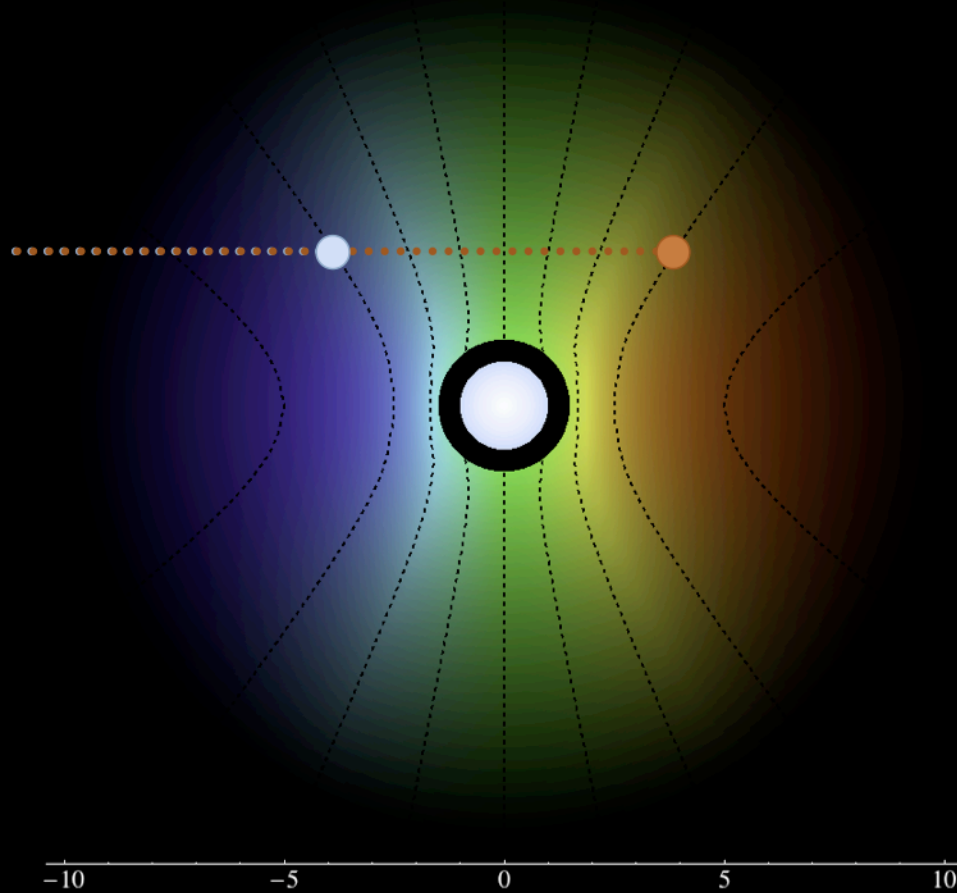


# Wind Profile Model



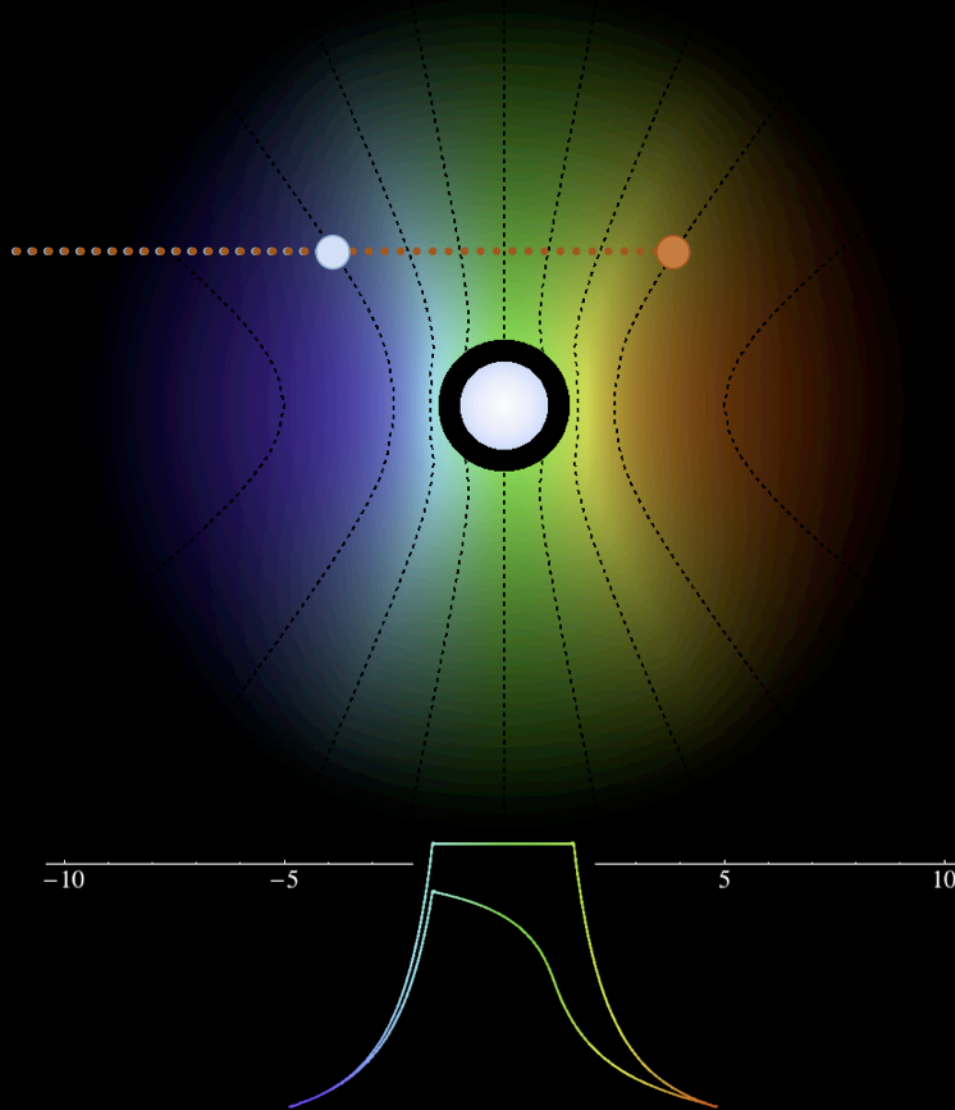
# Wind Profile Model

A



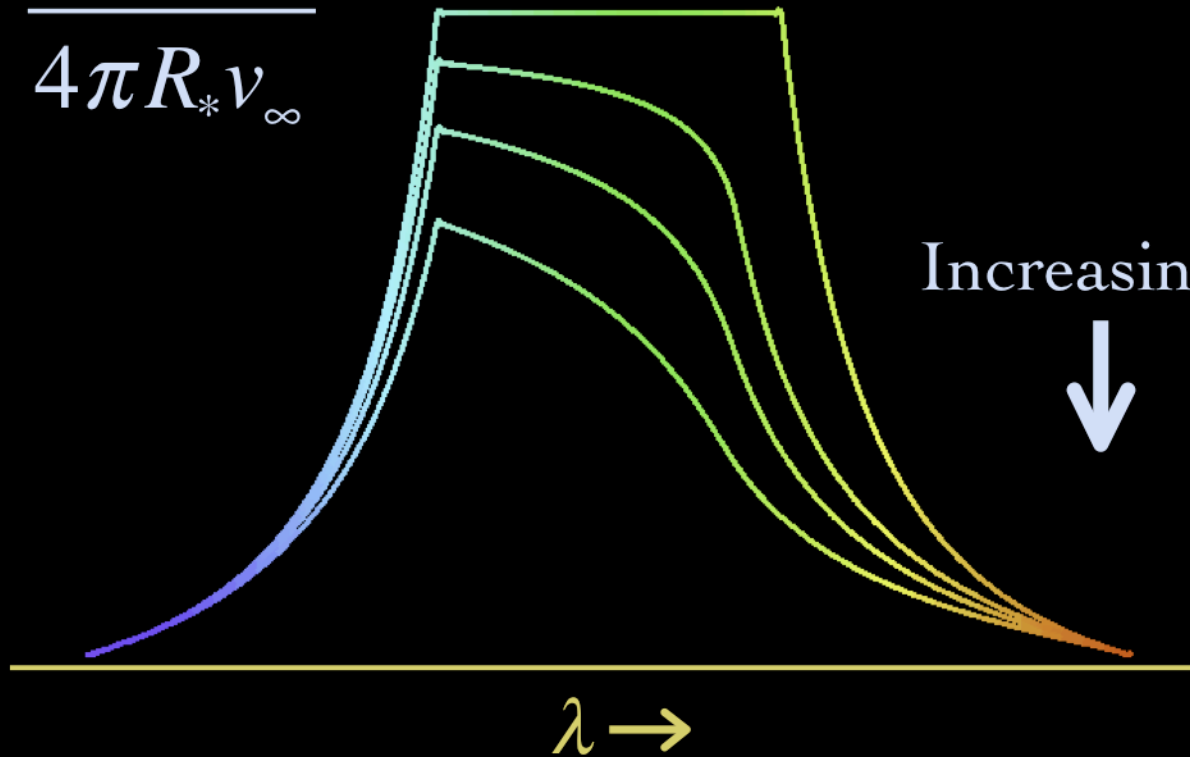
# Wind Profile Model

A



# Wind Profile Model

$$\tau_* = \frac{\kappa \dot{M}}{4\pi R_* v_\infty}$$



opacity of the cold  
wind component

wind mass-loss rate

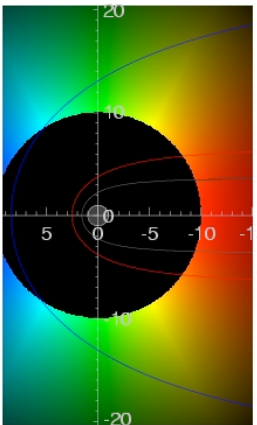
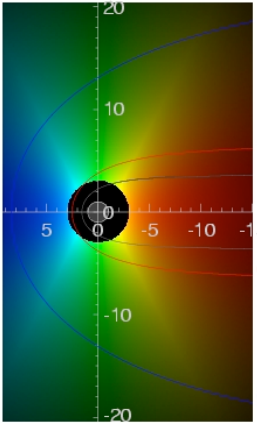
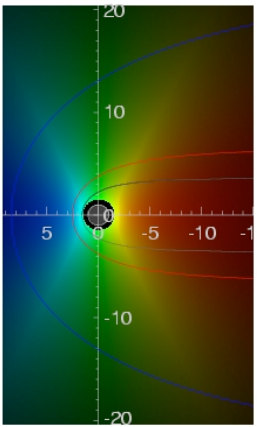
$$\dot{M} = 4\pi r^2 v \rho$$

$$\tau_* \equiv \frac{\kappa \dot{M}}{4\pi R_* v_\infty}$$

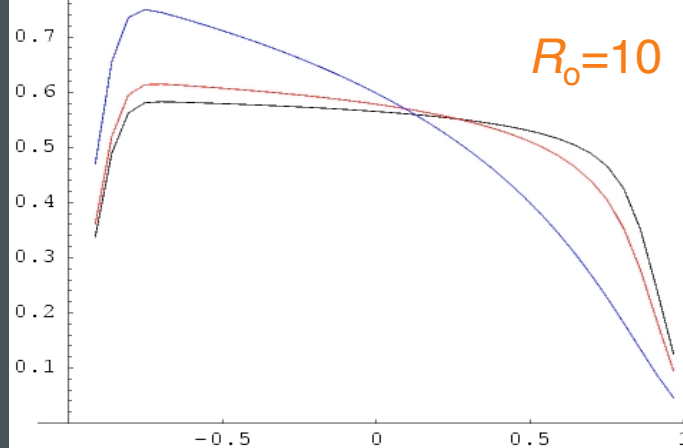
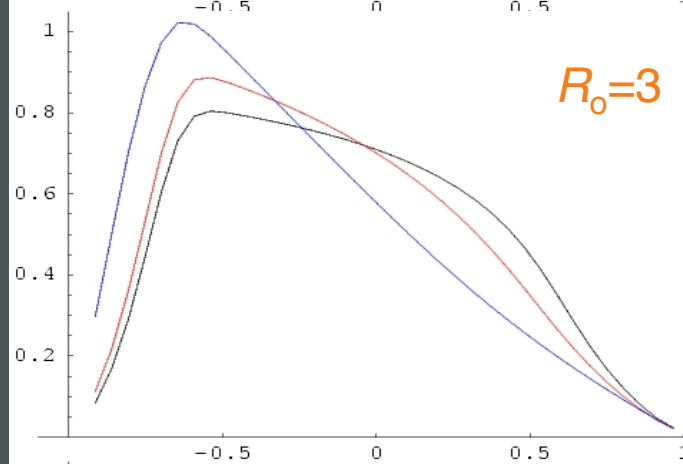
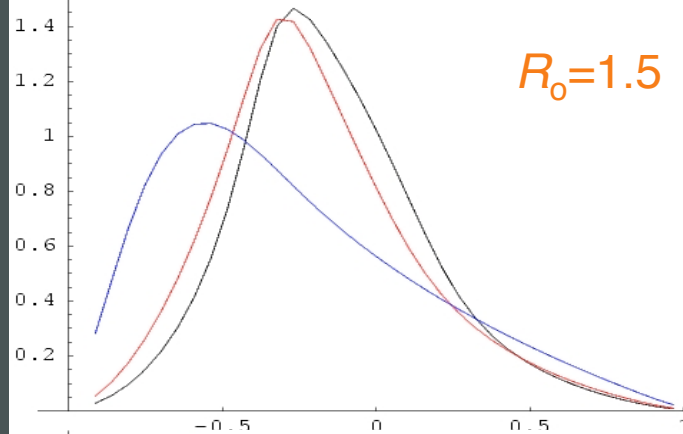
radius of the star

wind terminal velocity

$\tau=1$  contours



$\tau_* = 1, 2, 8$



## The basic wind-profile model

key parameters:  $R_0$  &  $\tau_*$

$$j \sim \rho^2 \text{ for } r/R_* > R_0, \\ = 0 \text{ otherwise}$$

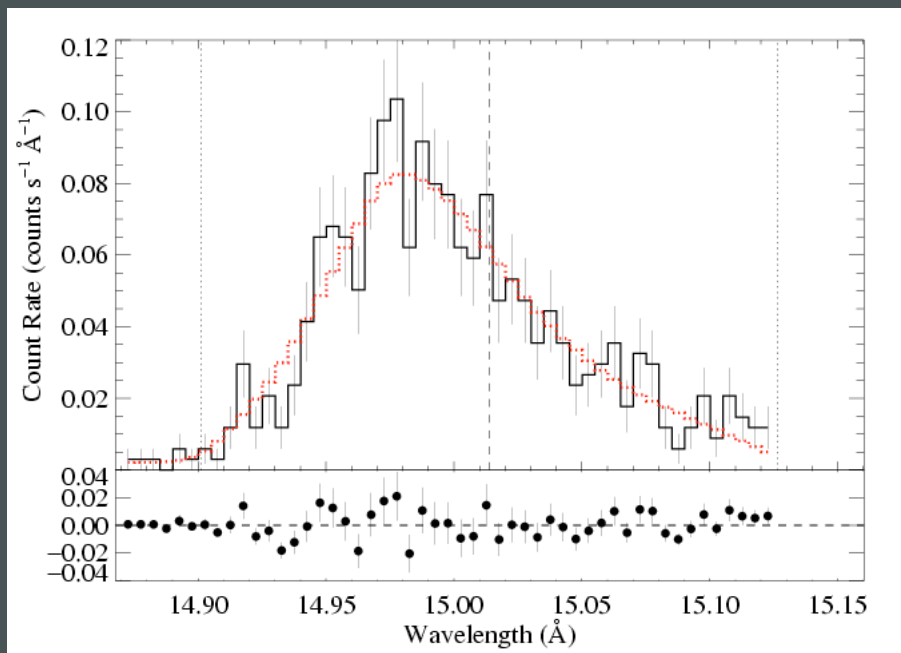
$$\tau = \tau_* \int_z^\infty \frac{R_* dz'}{r'^2 (1 - R_*/r')^\beta}$$

$$\tau_* \equiv \frac{\kappa \dot{M}}{4\pi R_* v_\infty}$$

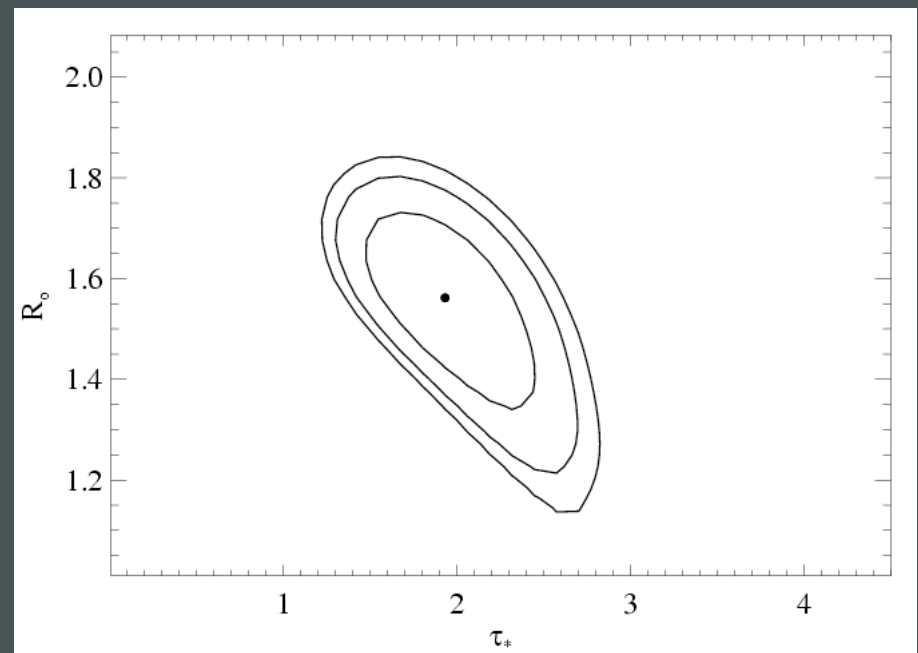


We fit these x-ray line profile models to each line in the *Chandra* data

And find a best-fit  $\tau_*$  and  $R_o$  & place confidence limits on these fitted parameter values

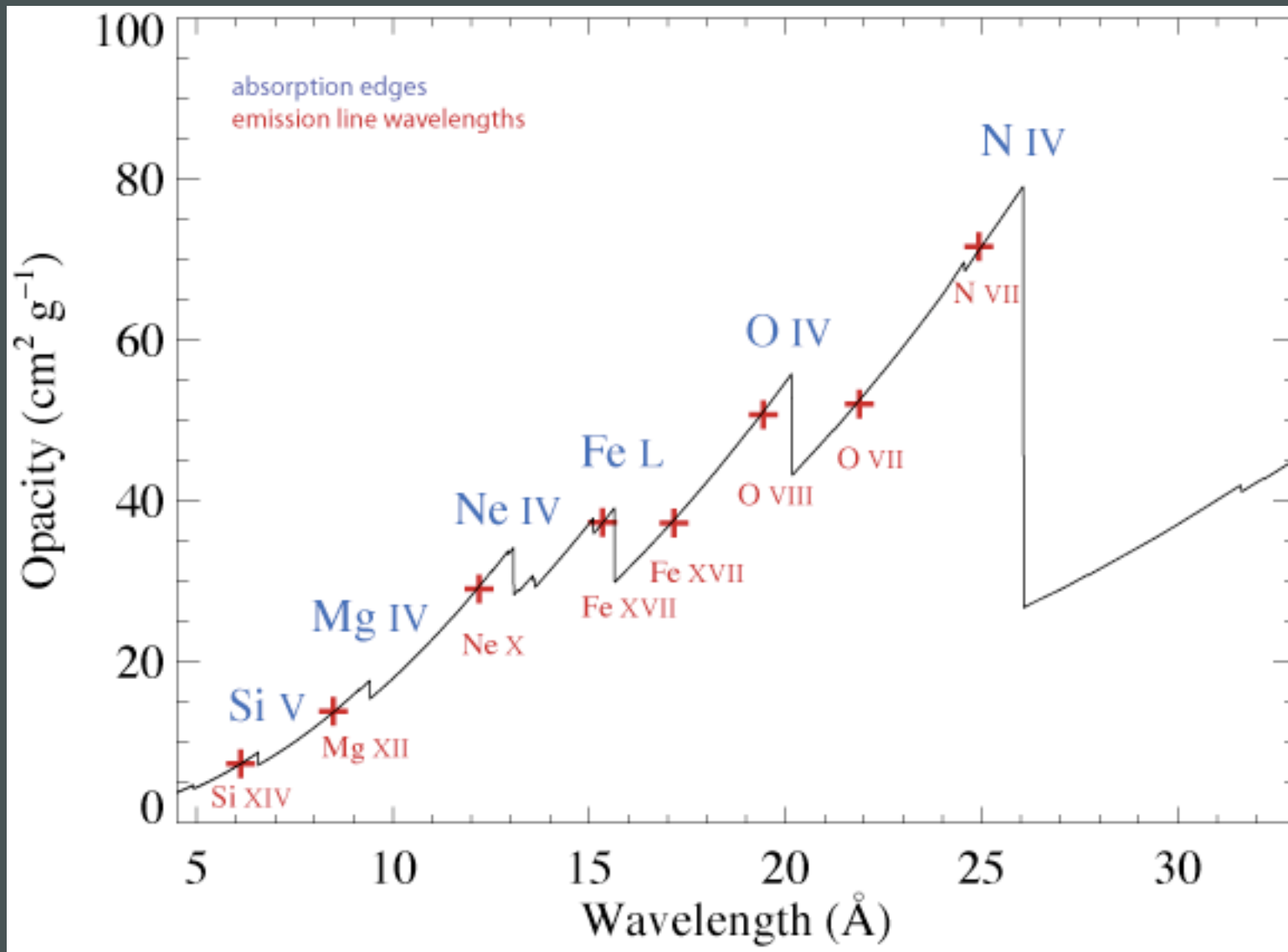


Fe XVII



68, 90, 95% confidence limits

# Wind opacity: photoelectric absorption



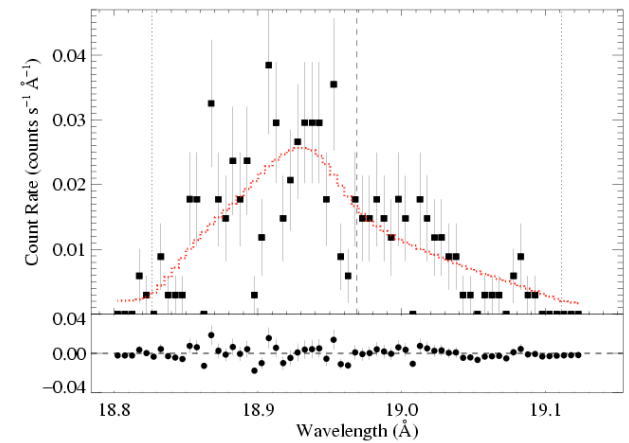
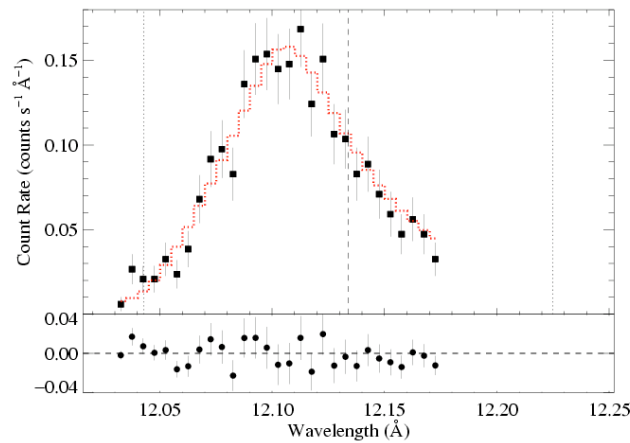
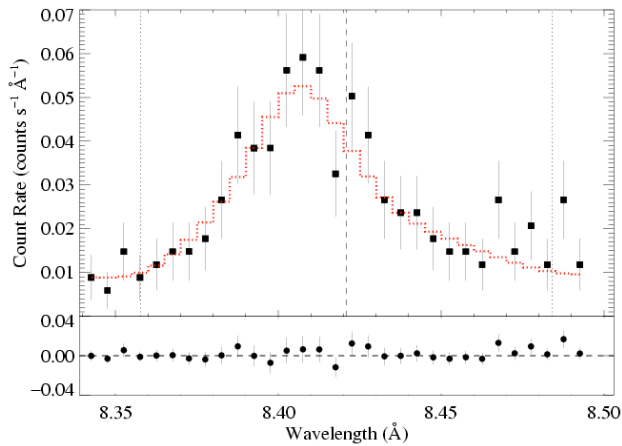
Abundances; ionization balance; atomic cross sections  
Verner & Yakovlev 1996

# $\zeta$ Pup: three emission lines

Mg Ly $\alpha$ : 8.42 Å

Ne Ly $\alpha$ : 12.13 Å

O Ly $\alpha$ : 18.97 Å



$$\tau_* = 1$$

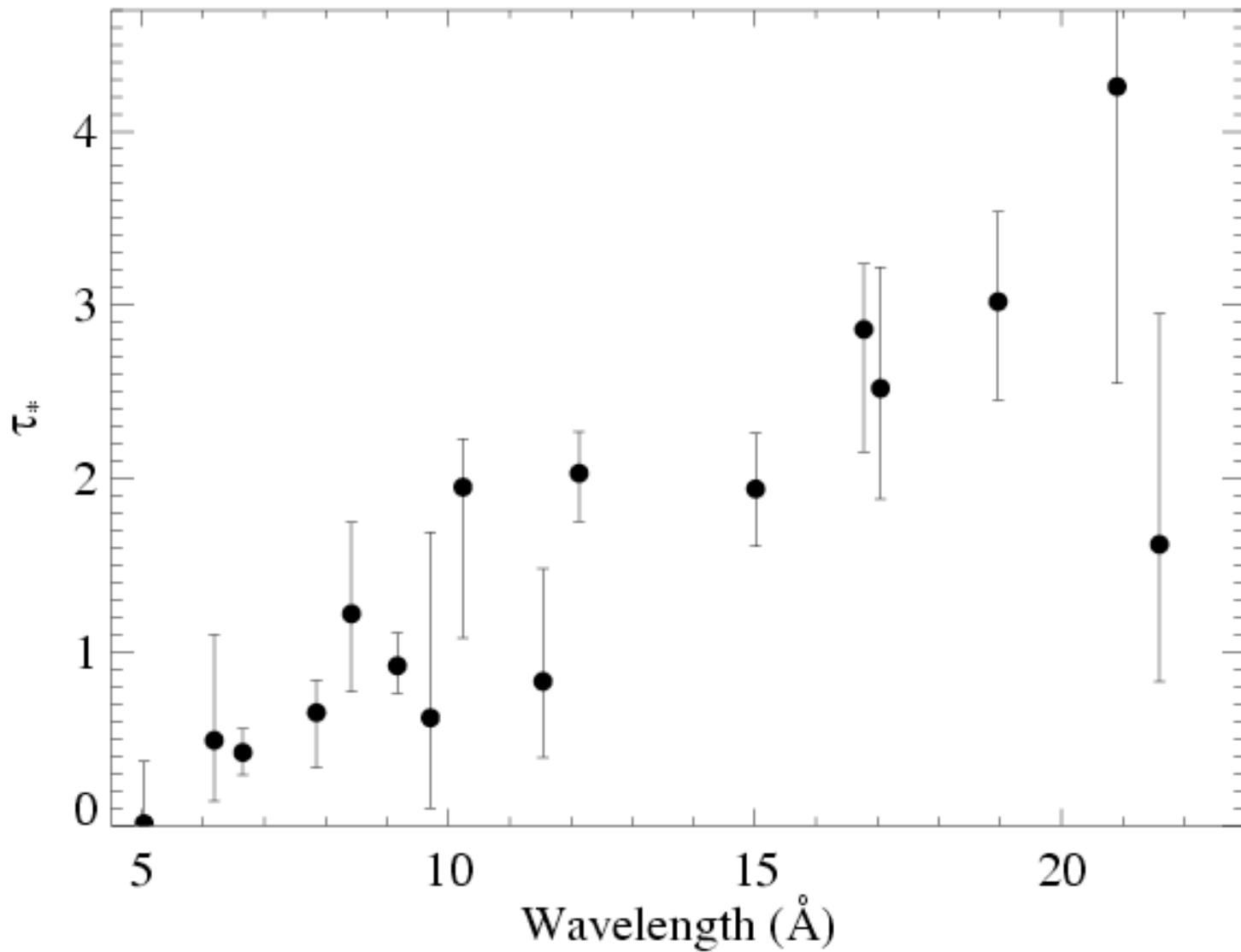
$$\tau_* = 2$$

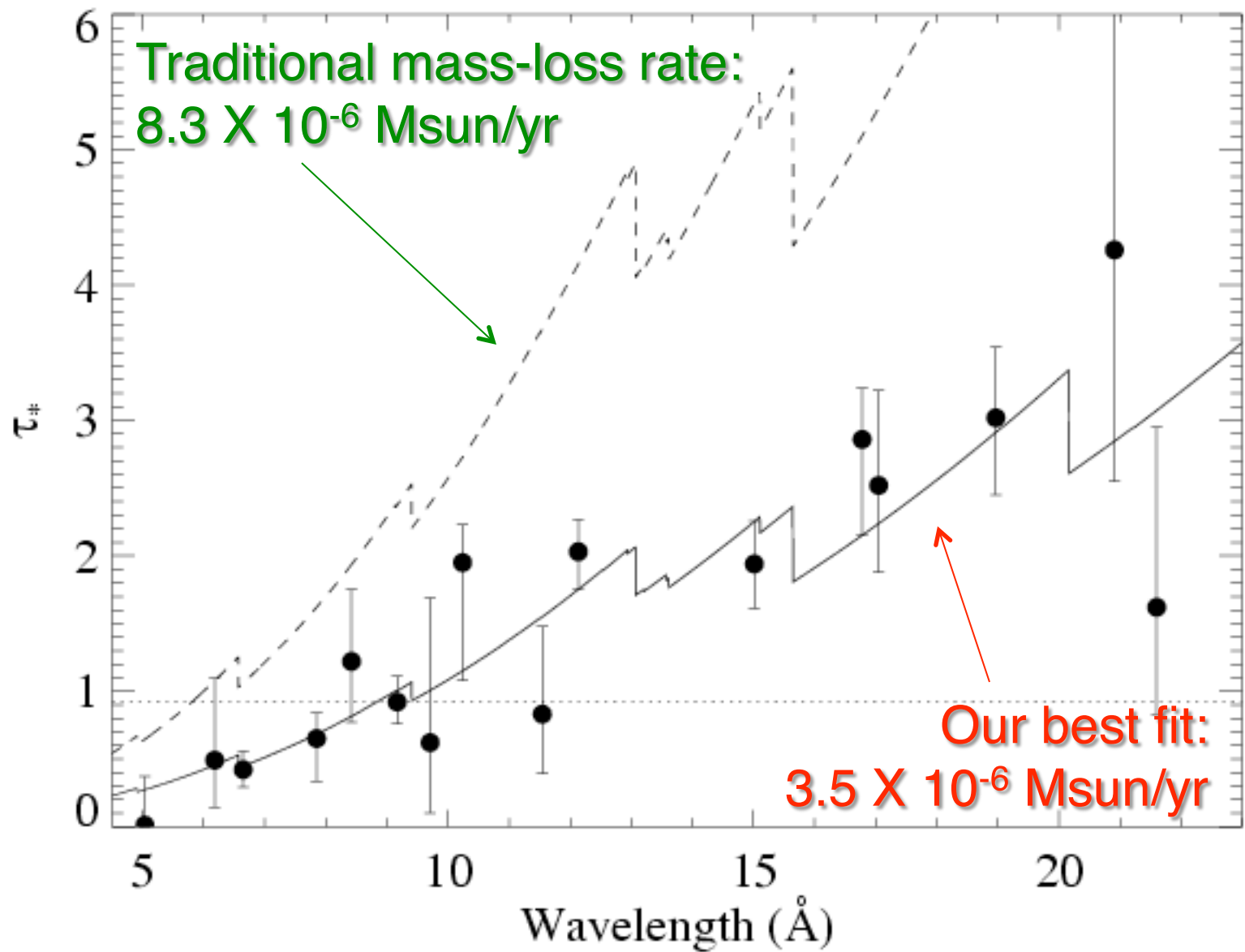
$$\tau_* = 3$$

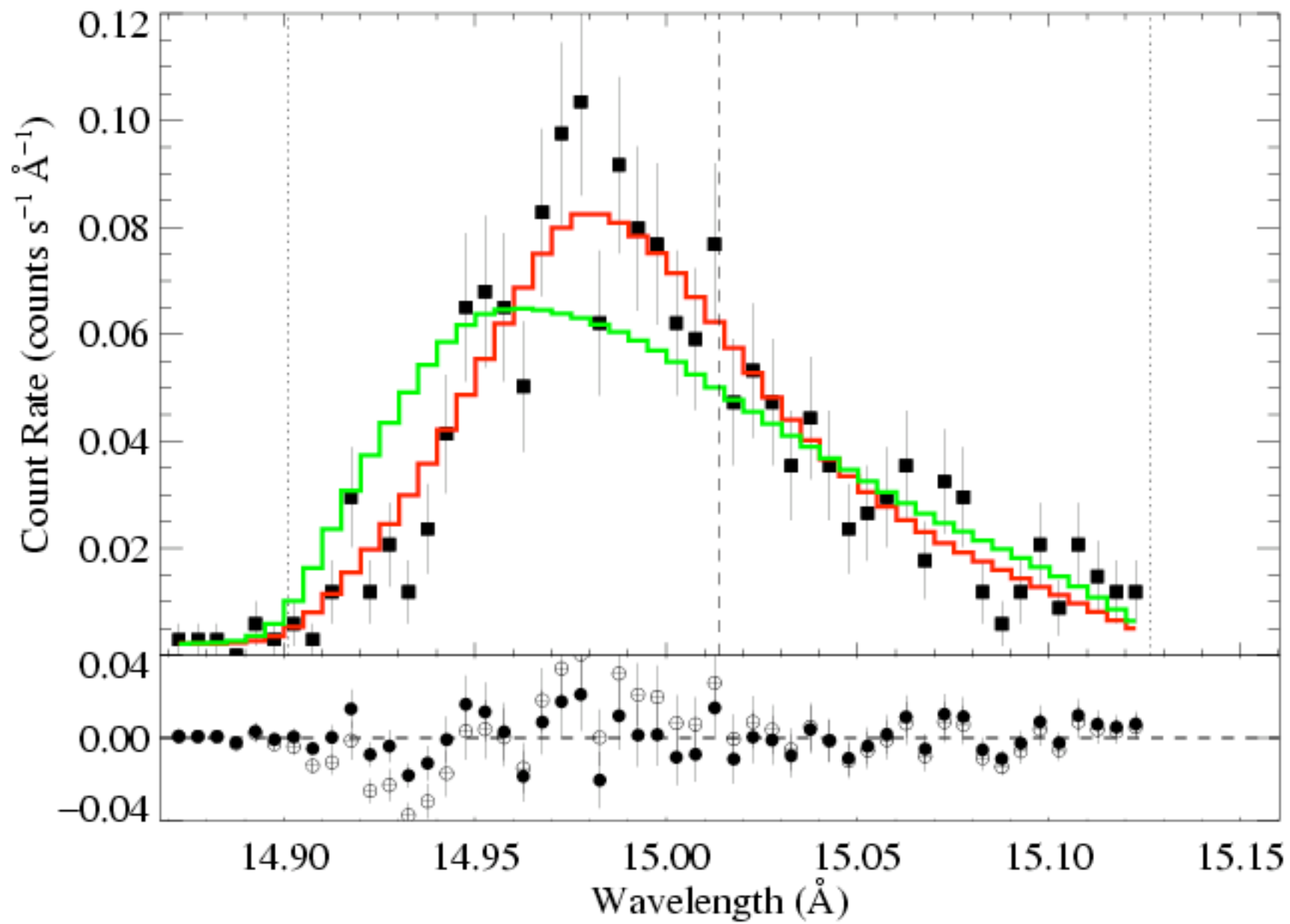
Recall:

$$\tau_* \equiv \frac{\kappa \dot{M}}{4\pi R_* v_\infty}$$

# Fits to 16 lines in the *Chandra* spectrum of $\zeta$ Pup

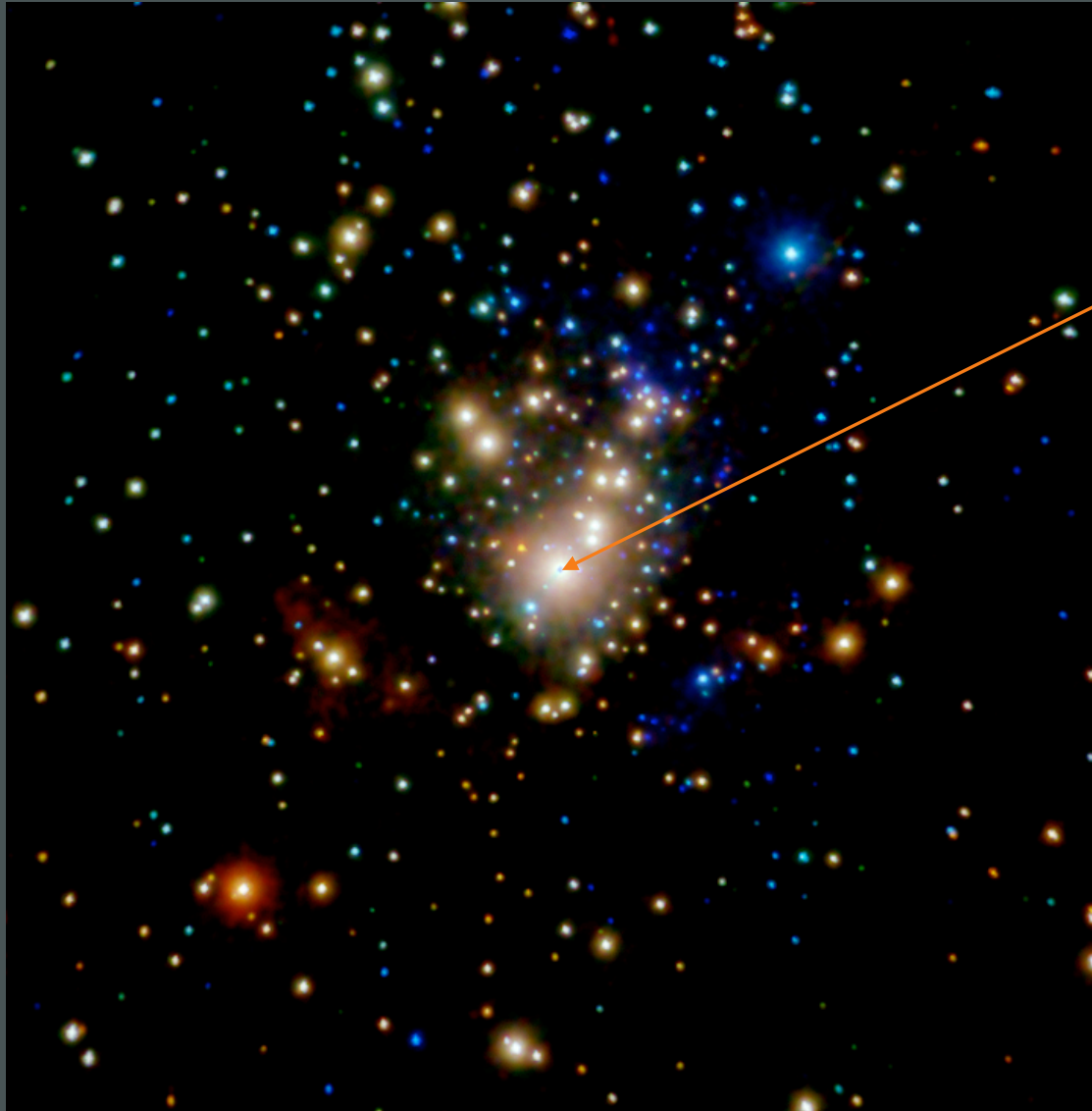






## Part 2: Magnetically Channeled Winds

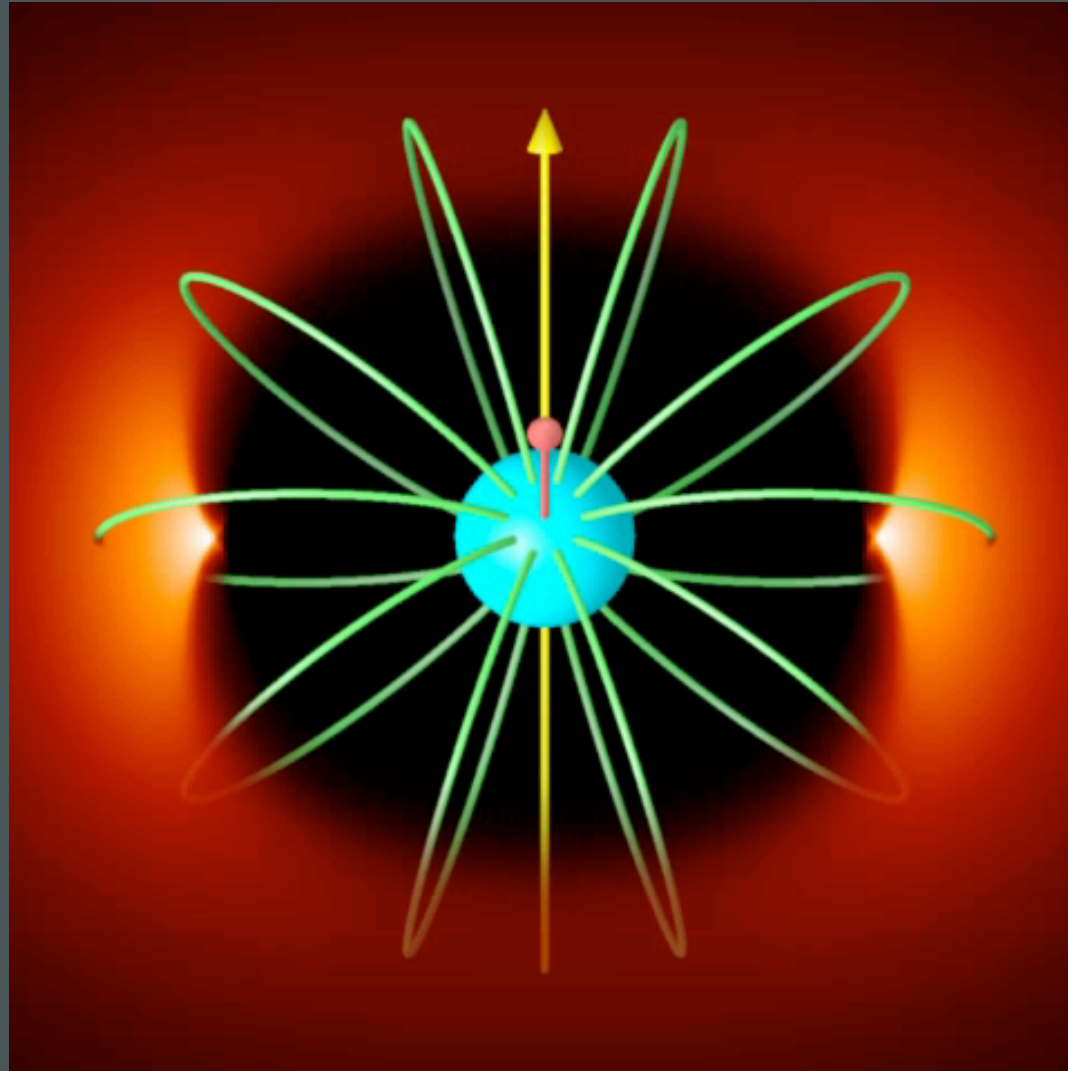
The central star in the Orion Nebula Cluster -  $\theta^1$  Ori C - is a source of strong and relatively hard x-rays



*Chandra* image: color coded by photon energy (red:  $<1$  keV; green 1 to 2 keV; blue  $> 2$  keV)



# Recently discovered dipole magnetic field of $> 1$ kG : Zeeman Doppler spectroscopy (Wade et al. 2006)



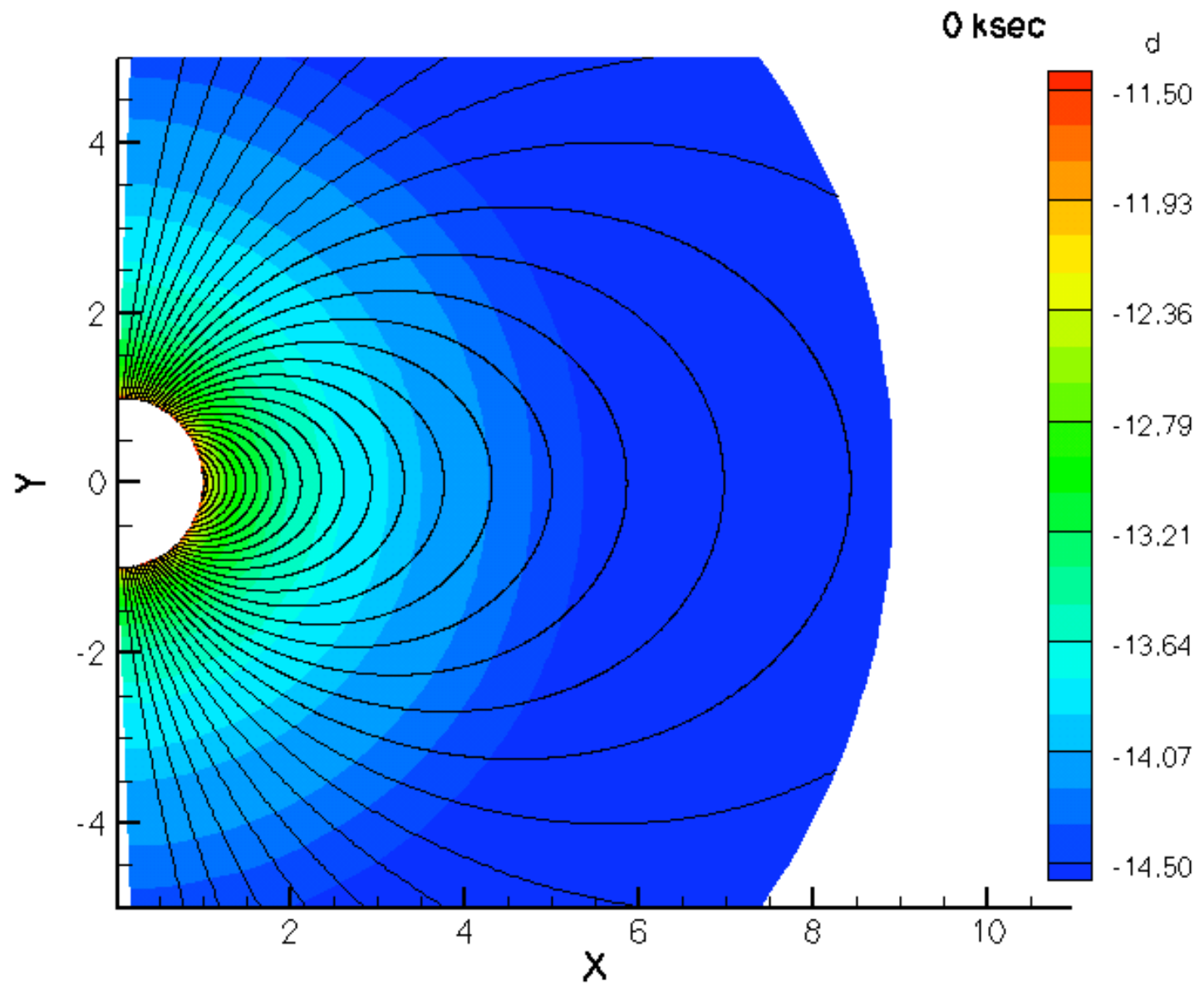
Simulation/visualization courtesy R. Townsend

Movie available at [astro.swarthmore.edu/~cohen/presentations/apip09/rrm-o25-i75-b60-redt.avi](http://astro.swarthmore.edu/~cohen/presentations/apip09/rrm-o25-i75-b60-redt.avi)

This field confines and channels the stellar wind out to several stellar radii –

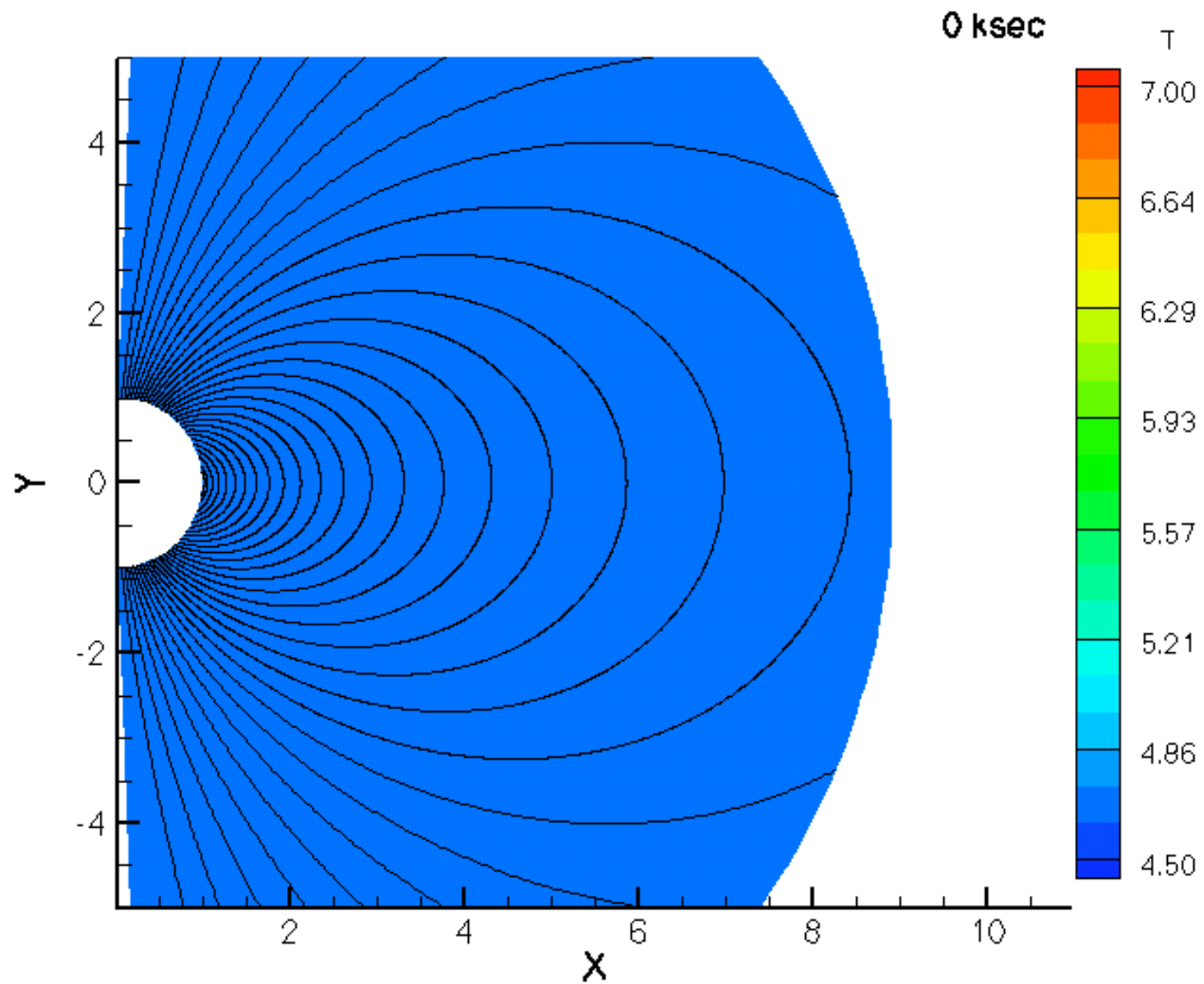
what is the effect?

...MHD simulations (ZEUS, 2-D)



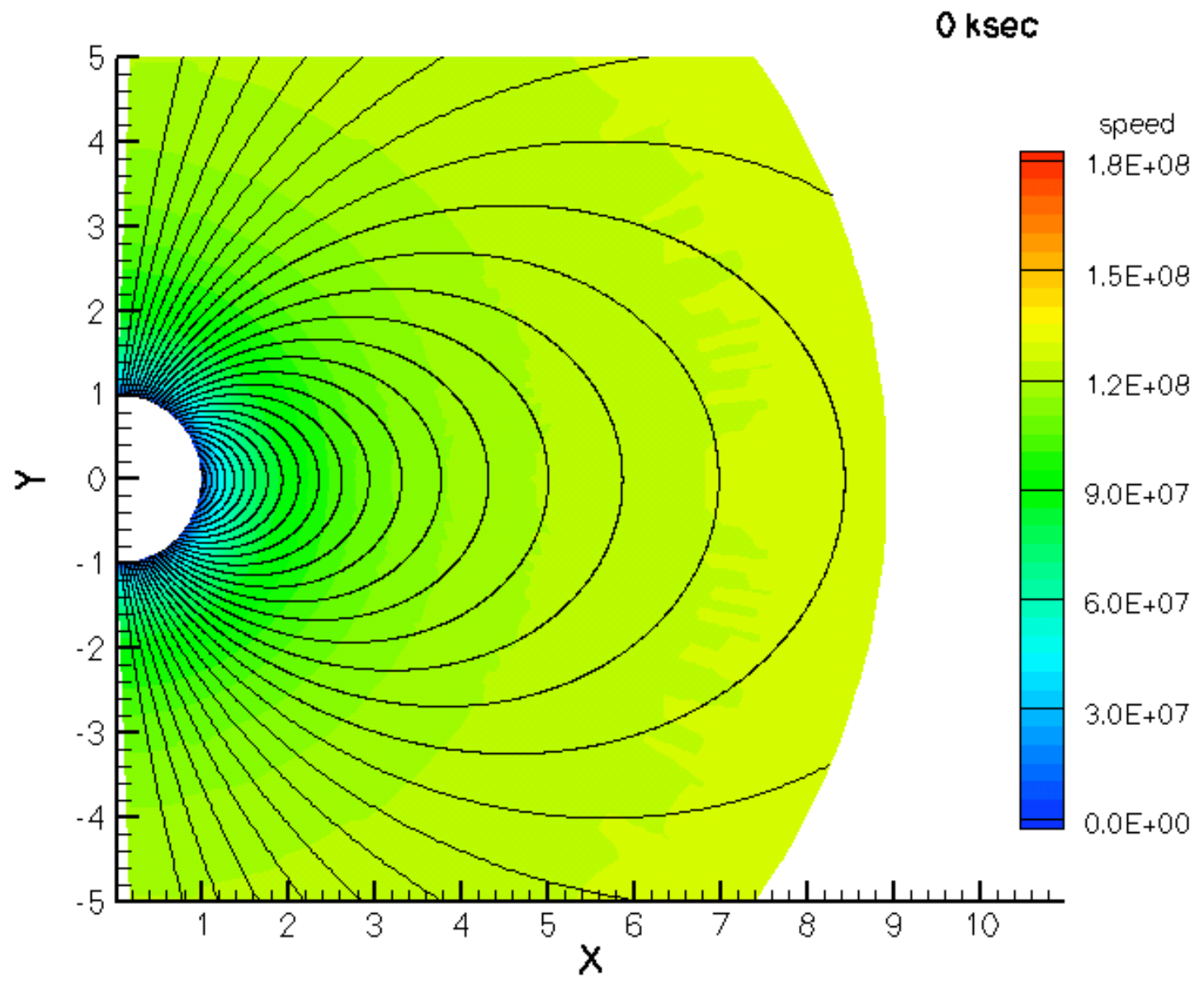
Simulation/visualization courtesy A. ud-Doula

Movie available at [astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-logd.avi](http://astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-logd.avi)



Simulation/visualization courtesy A. ud-Doula

Movie available at [astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-logT.avi](http://astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-logT.avi)

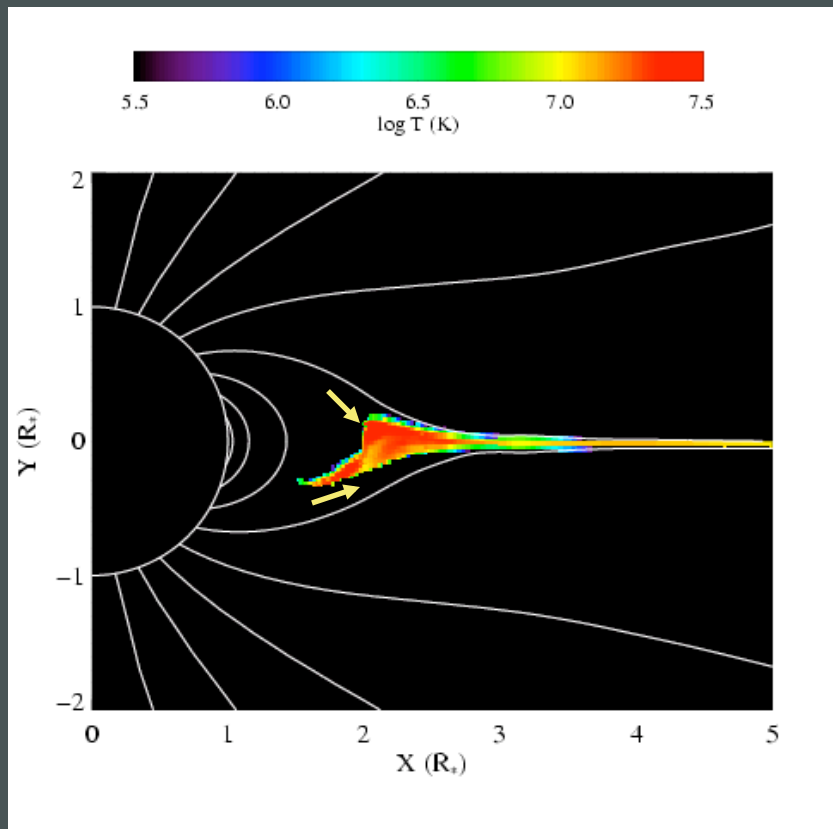


Simulation/visualization courtesy A. ud-Doula

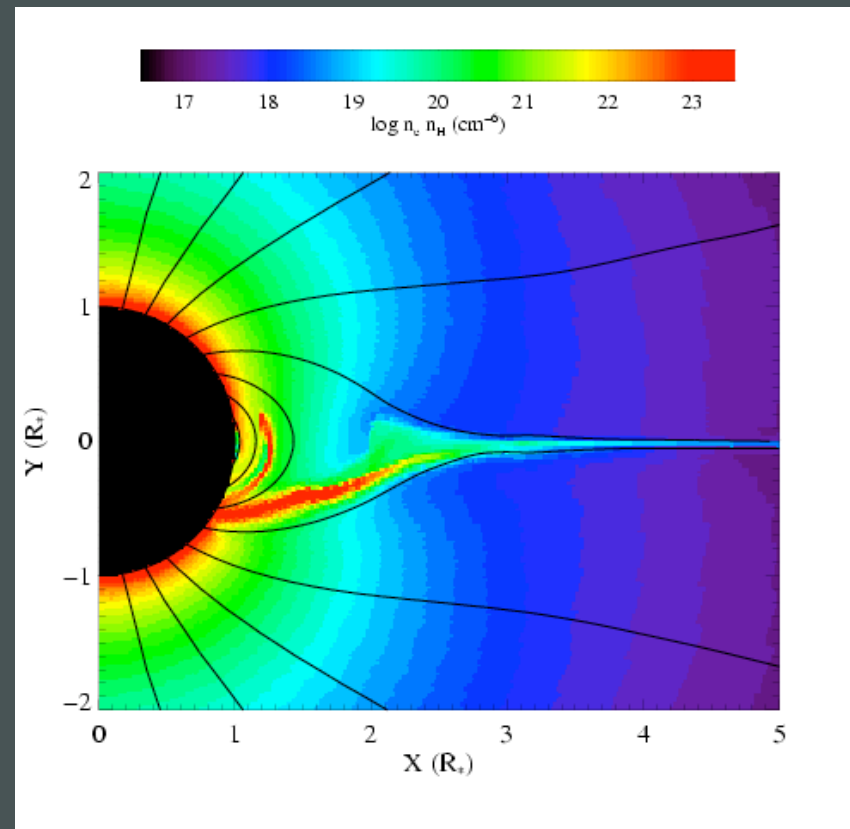
Movie available at [astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-speed.avi](http://astro.swarthmore.edu/~cohen/presentations/apip09/t1oc-lowvinf-speed.avi)

# MHD simulations of magnetically channeled wind

temperature



emission measure



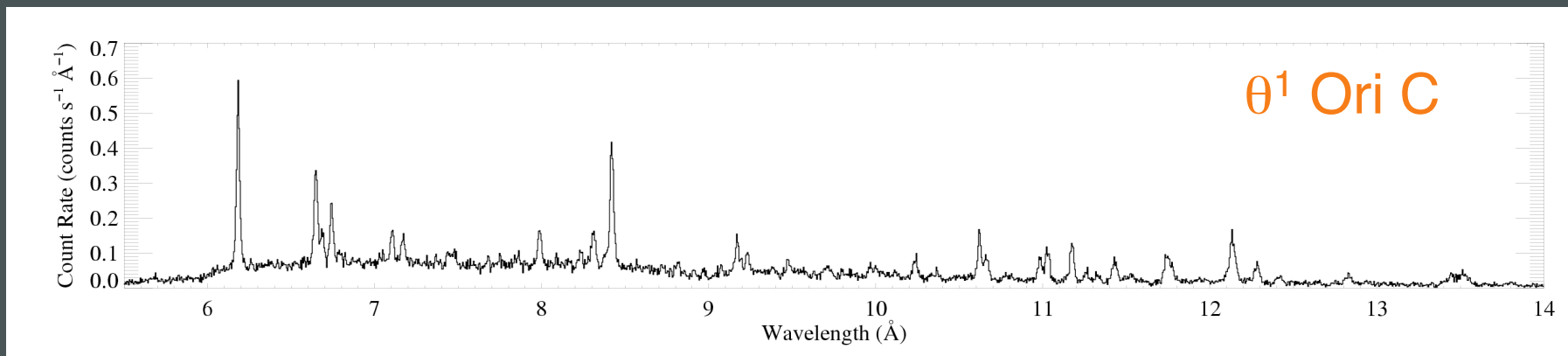
simulations by A. ud-Doula; Gagné et al. (2005)

Channeled collision is close to head-on –  
>1000 km s<sup>-1</sup> : T > 10<sup>7</sup> K

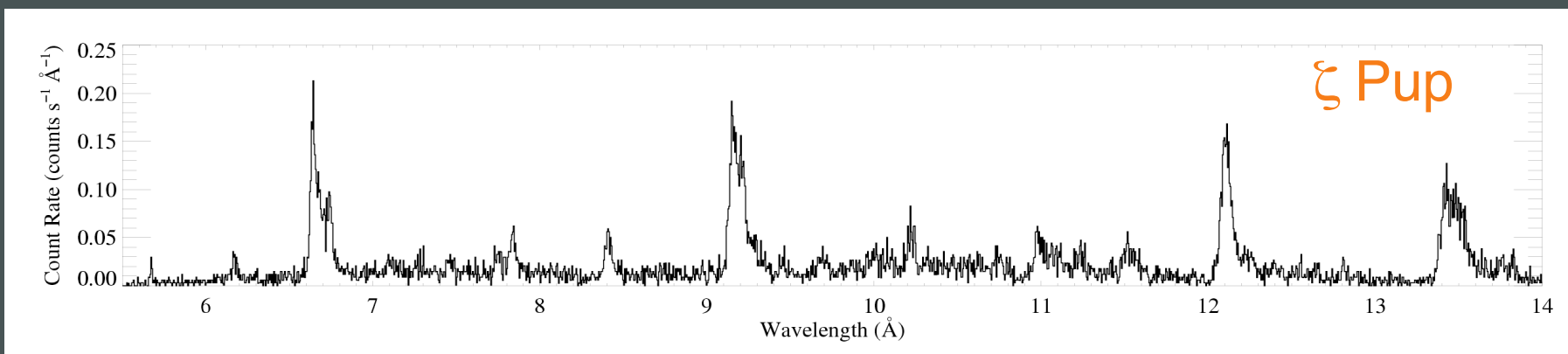
## Predictions:

1. Shocks are strong – head-on – and so plasma is hotter;
2. Hot plasma is moving much slower (confinement);
3. Rotational modulation of X-ray flux;
4. Hot plasma is  $\sim 1 R_*$  above the surface.

## Chandra grating spectra



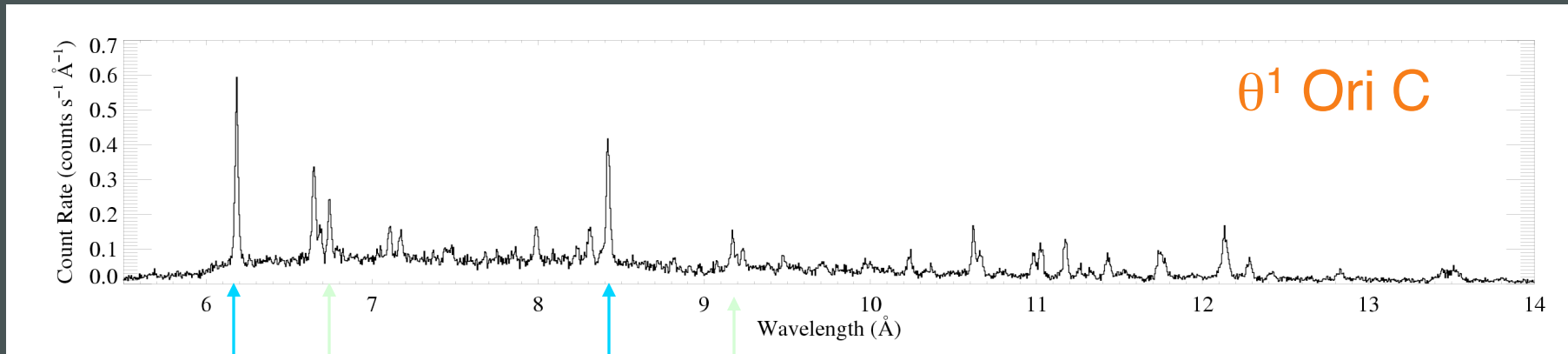
$\theta^1$  Ori C: hotter plasma, narrower emission lines



$\zeta$  Pup: cooler plasma, broad emission lines



# H-like/He-like ratio is temperature sensitive

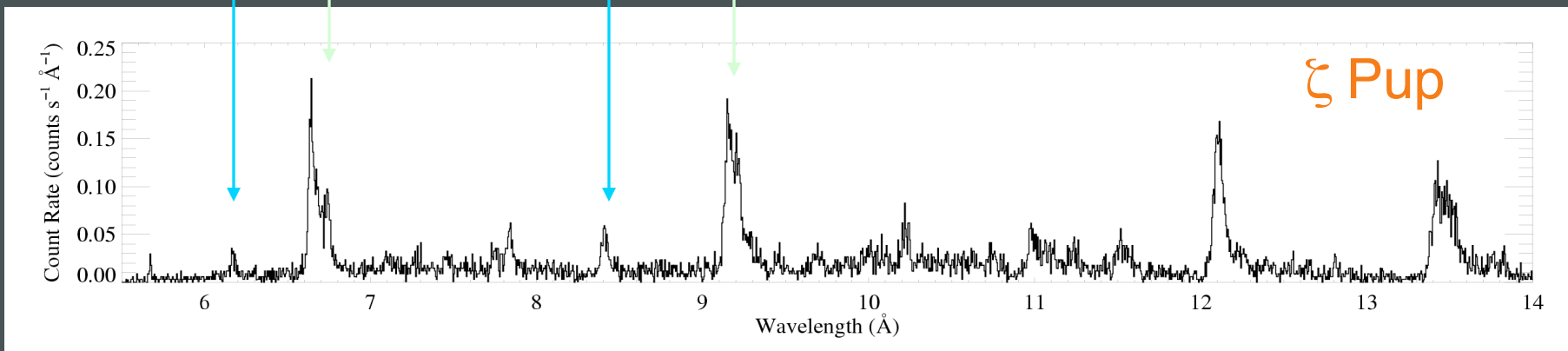


Si XIV

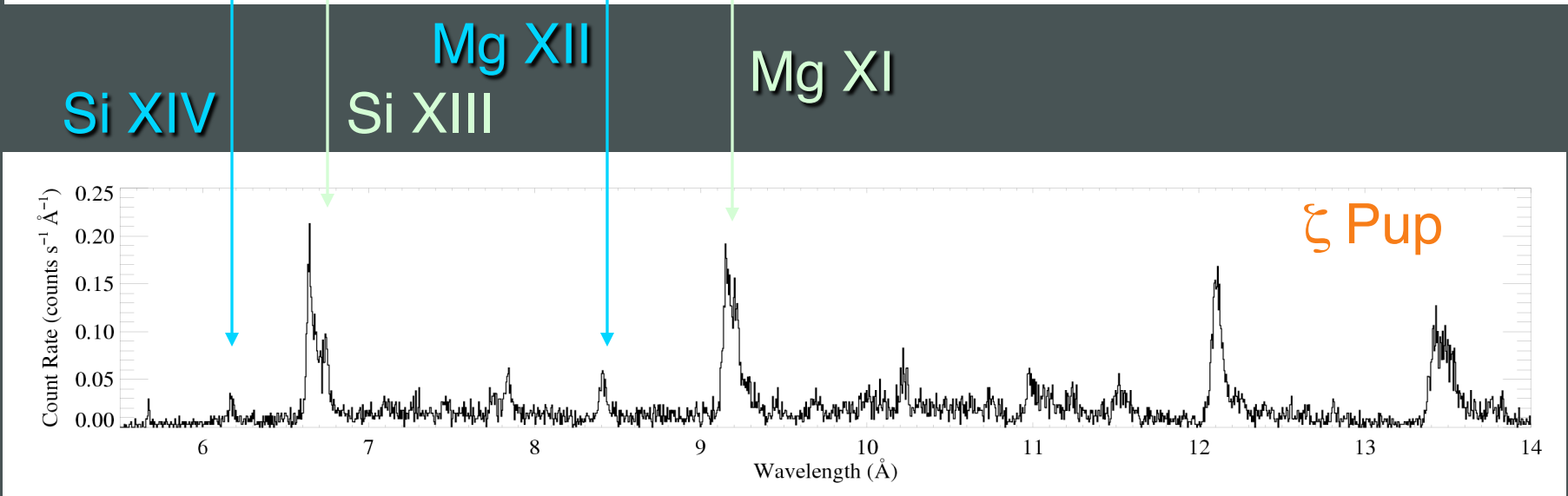
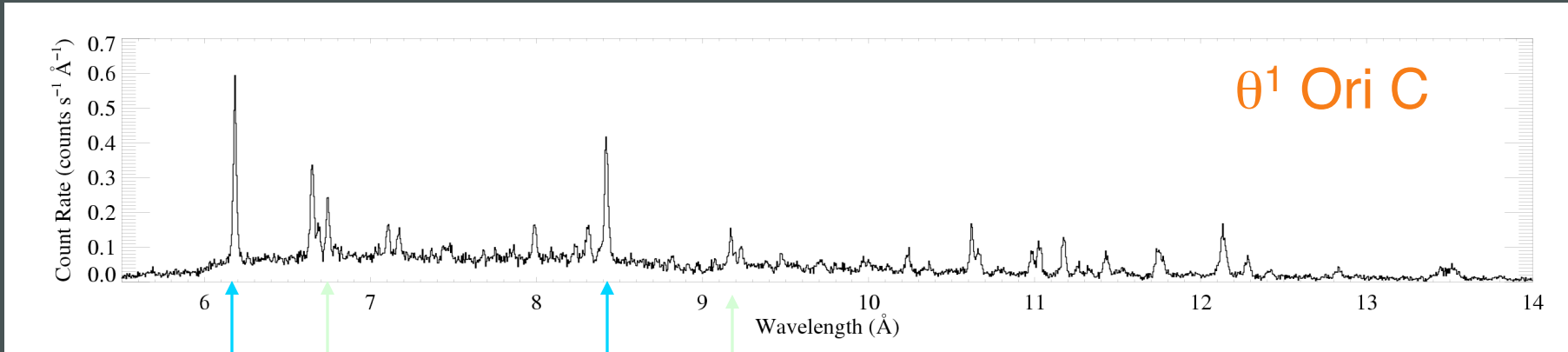
Mg XII

Si XIII

Mg XI

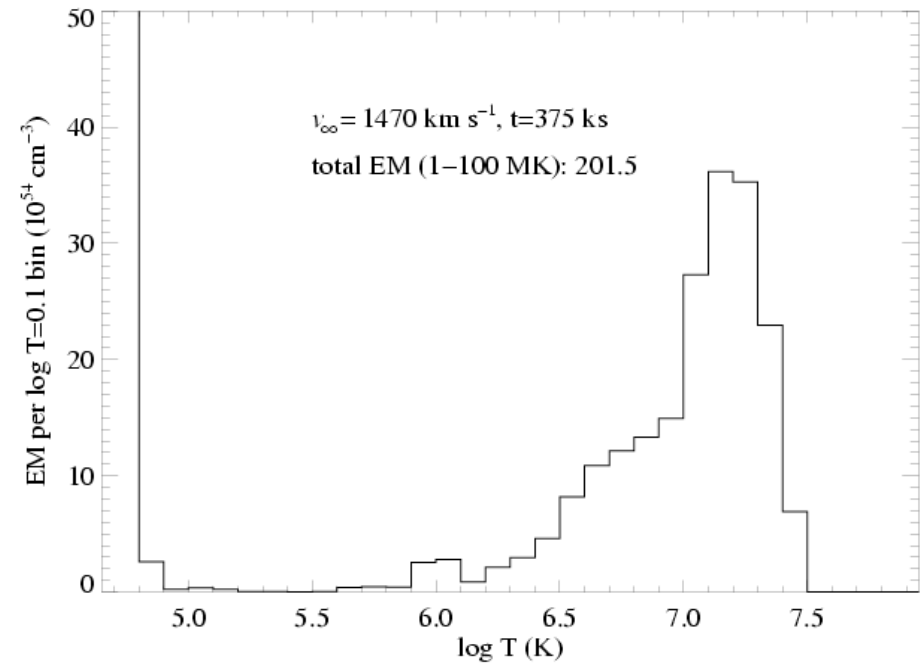
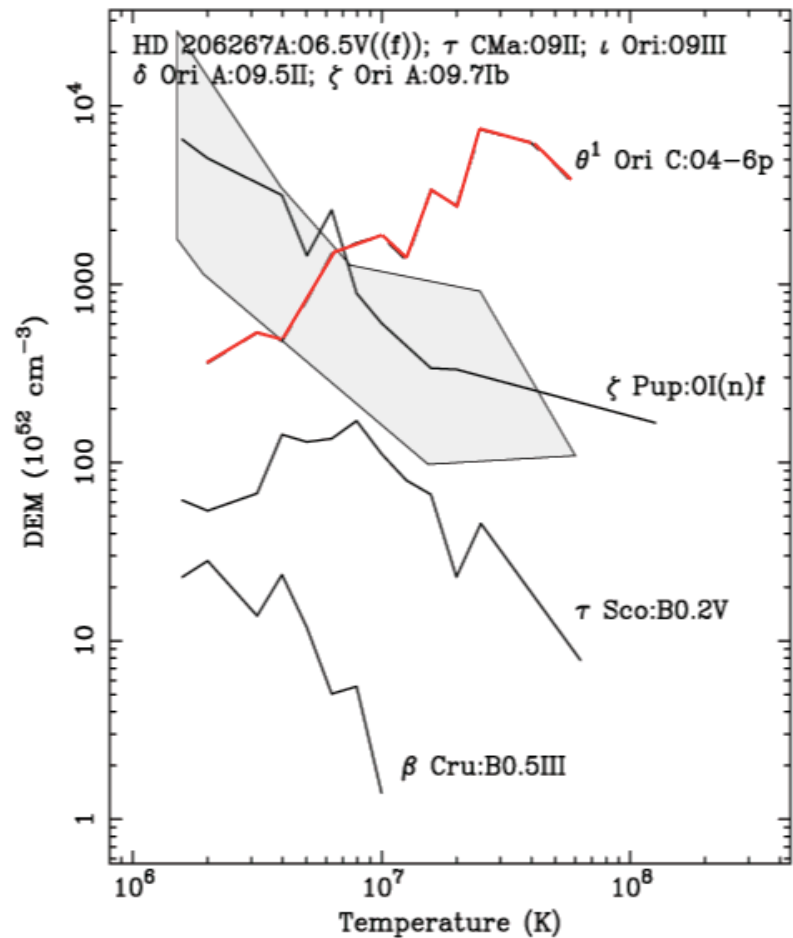


# The magnetic O star – $\theta^1$ Ori C – is hotter



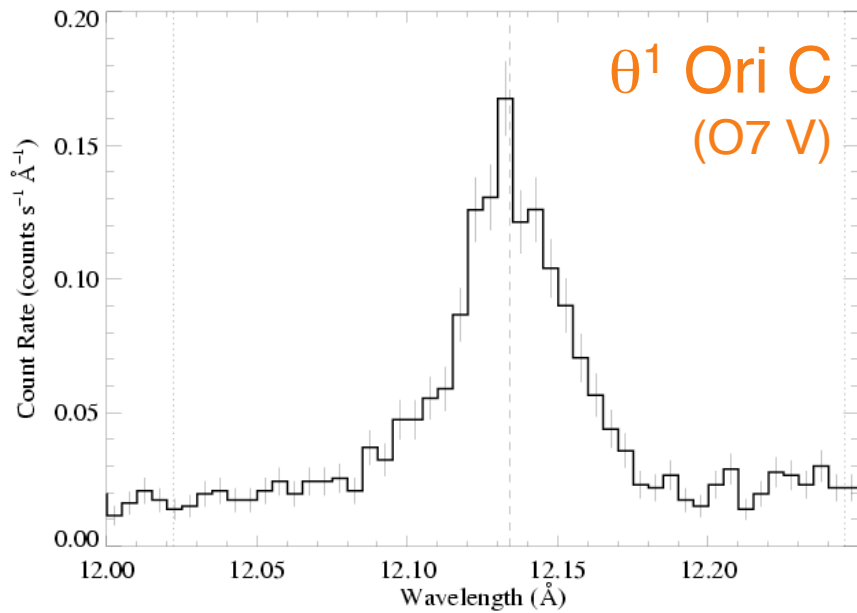
# Differential emission measure

(temperature distribution)

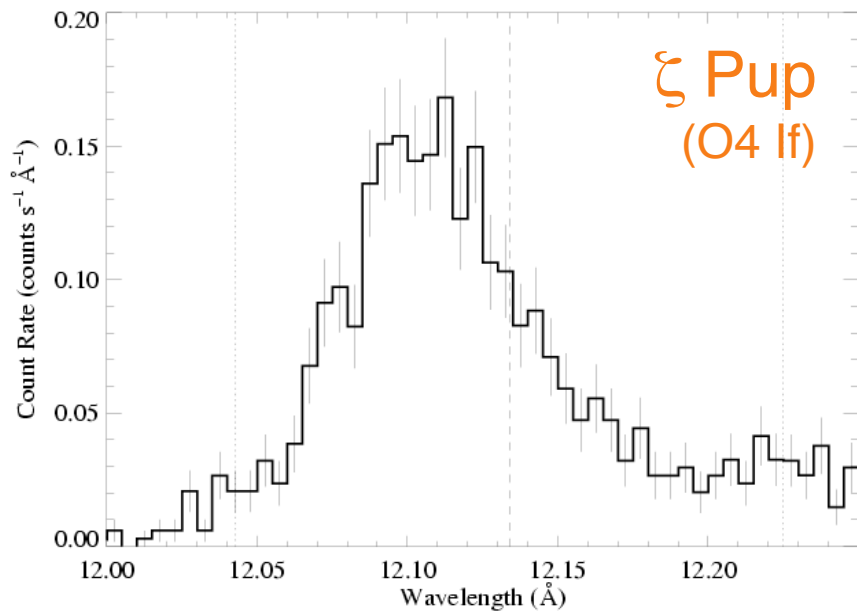


MHD simulation of  $\theta^1$  Ori C  
reproduces the observed  
differential emission measure

Wojdowski & Schulz (2005)

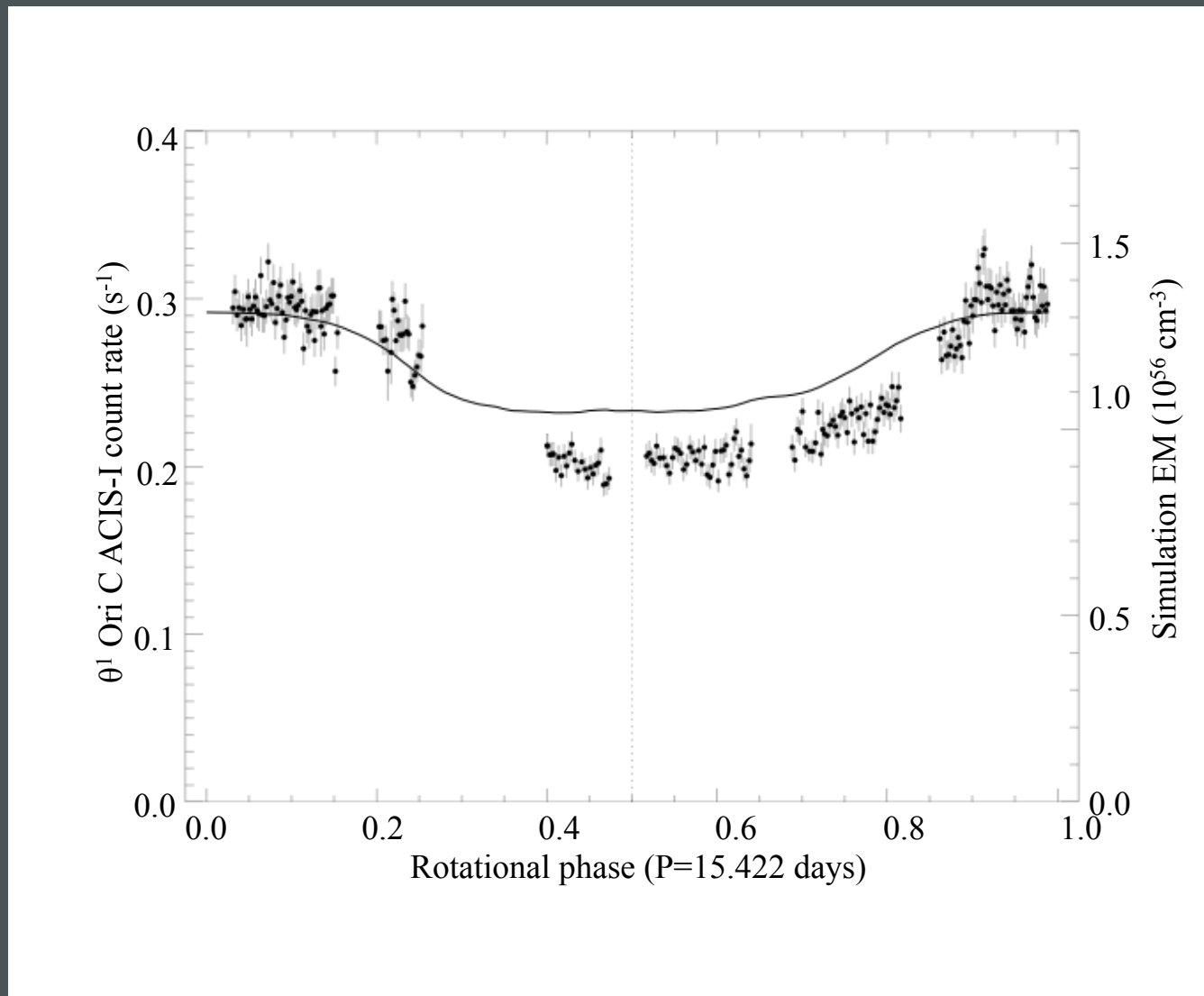


1000  $\text{km s}^{-1}$



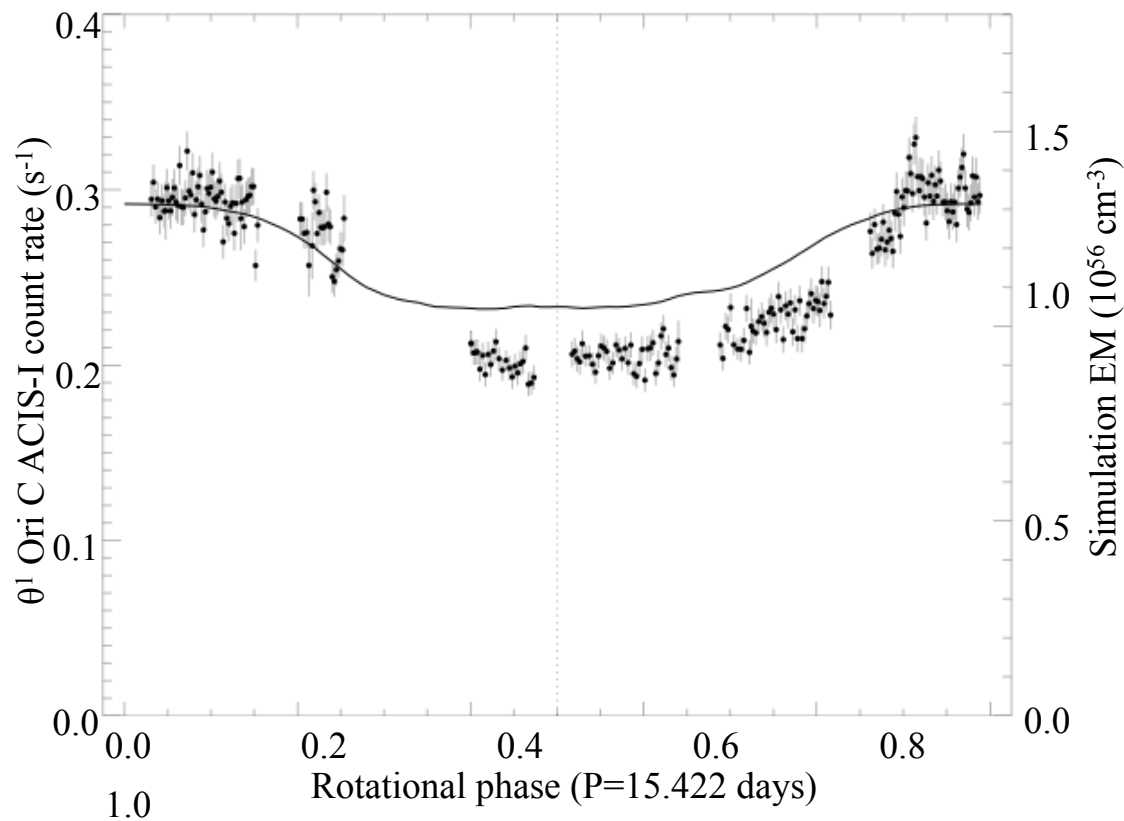
Emission lines *are* significantly narrower in the magnetic massive star's x-ray spectrum

## Chandra broadband count rate vs. rotational phase



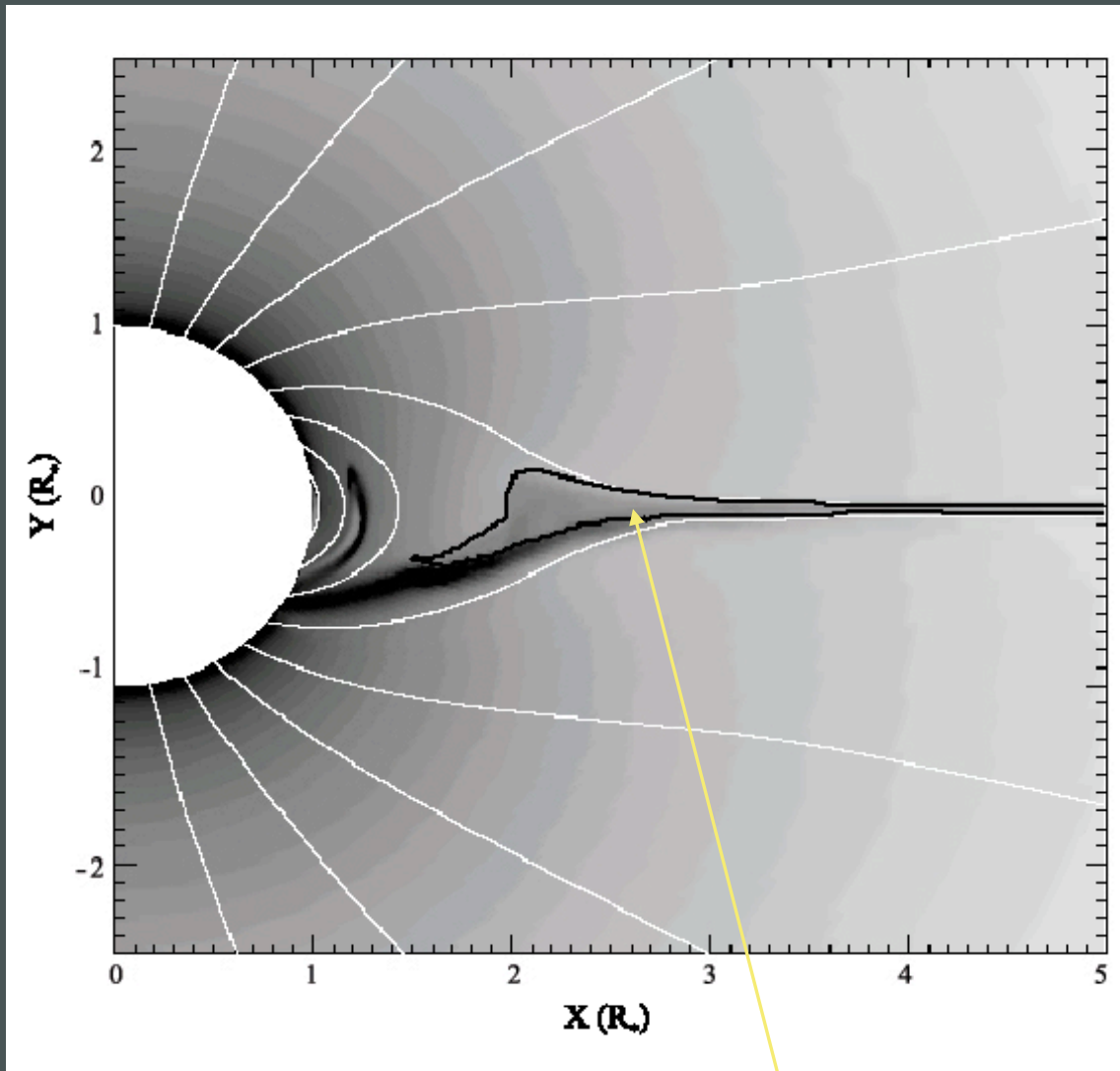
Model from MHD simulation

## The star itself occults the hot plasma torus



The closer the hot plasma is to the star, the deeper the dip in the x-ray light curve

## Emission measure



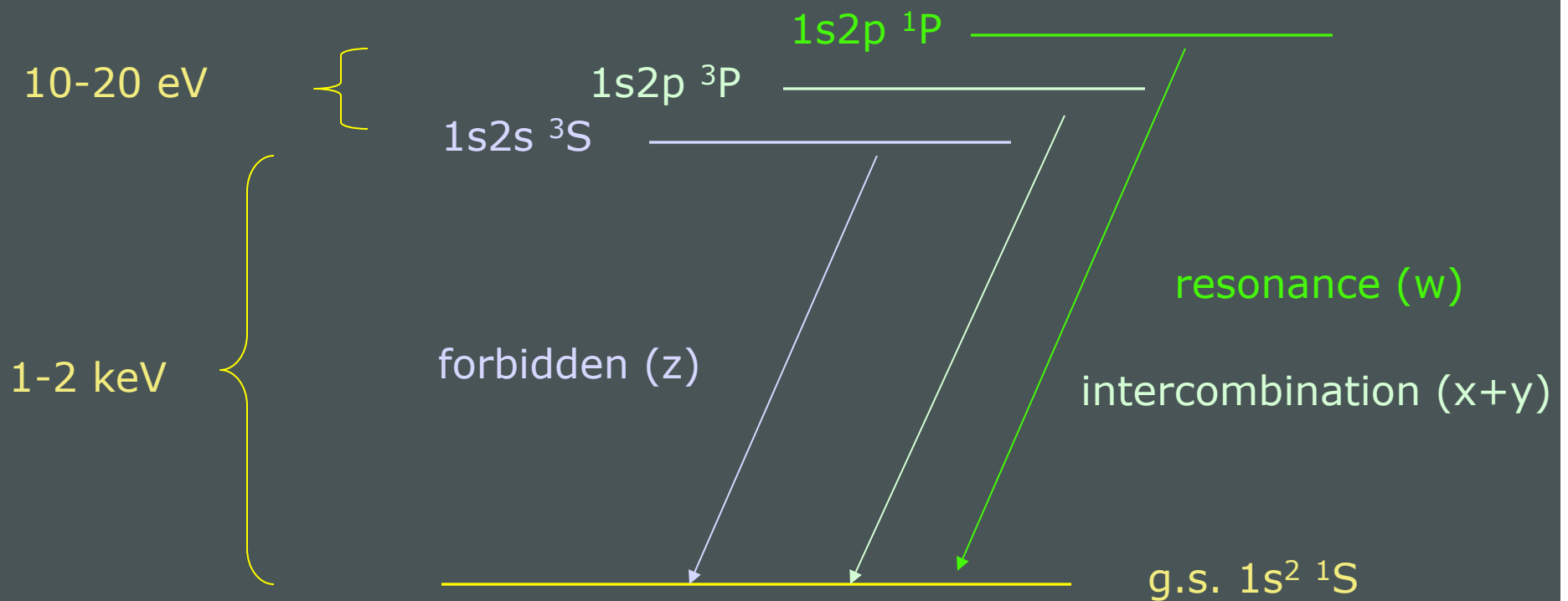
contour encloses  $T > 10^6$  K

Helium-like species' forbidden-to-intercombination line ratios –  $z/(x+y)$  – provide information about the *location* of the hot plasma

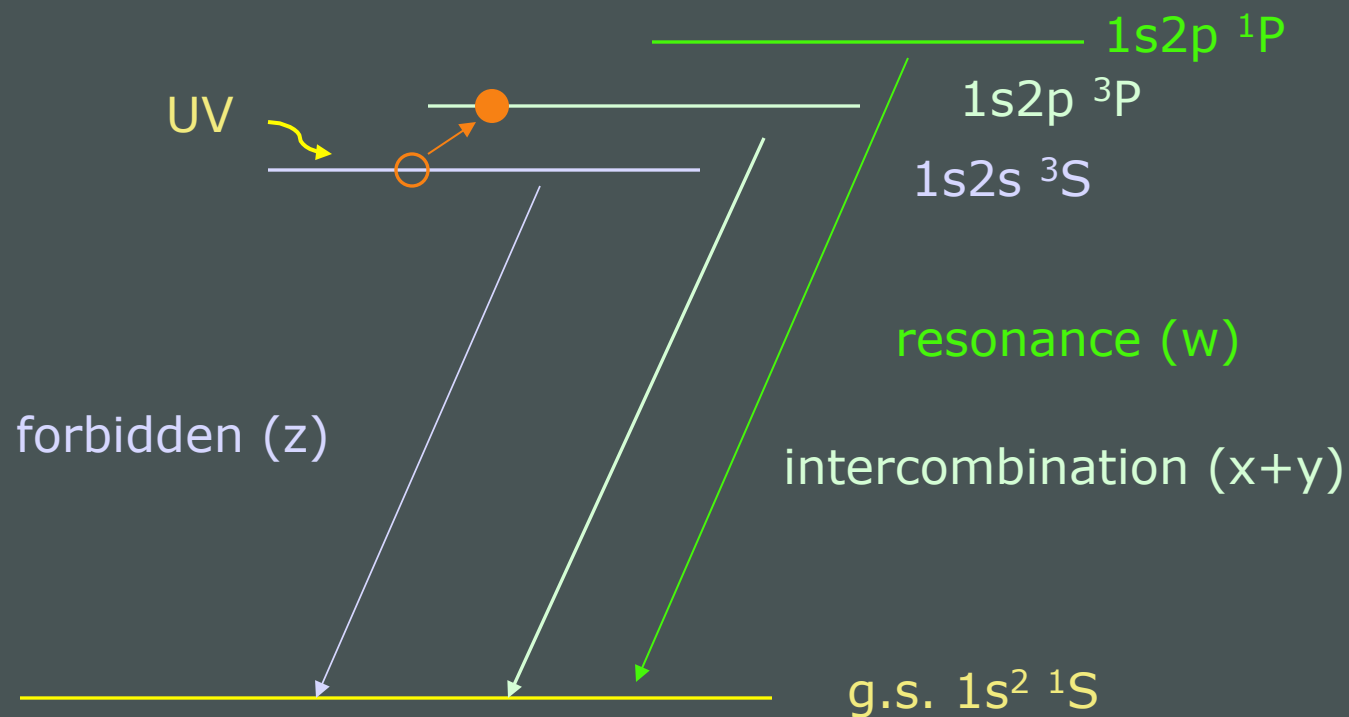
...not the *density*, as is usually the case.



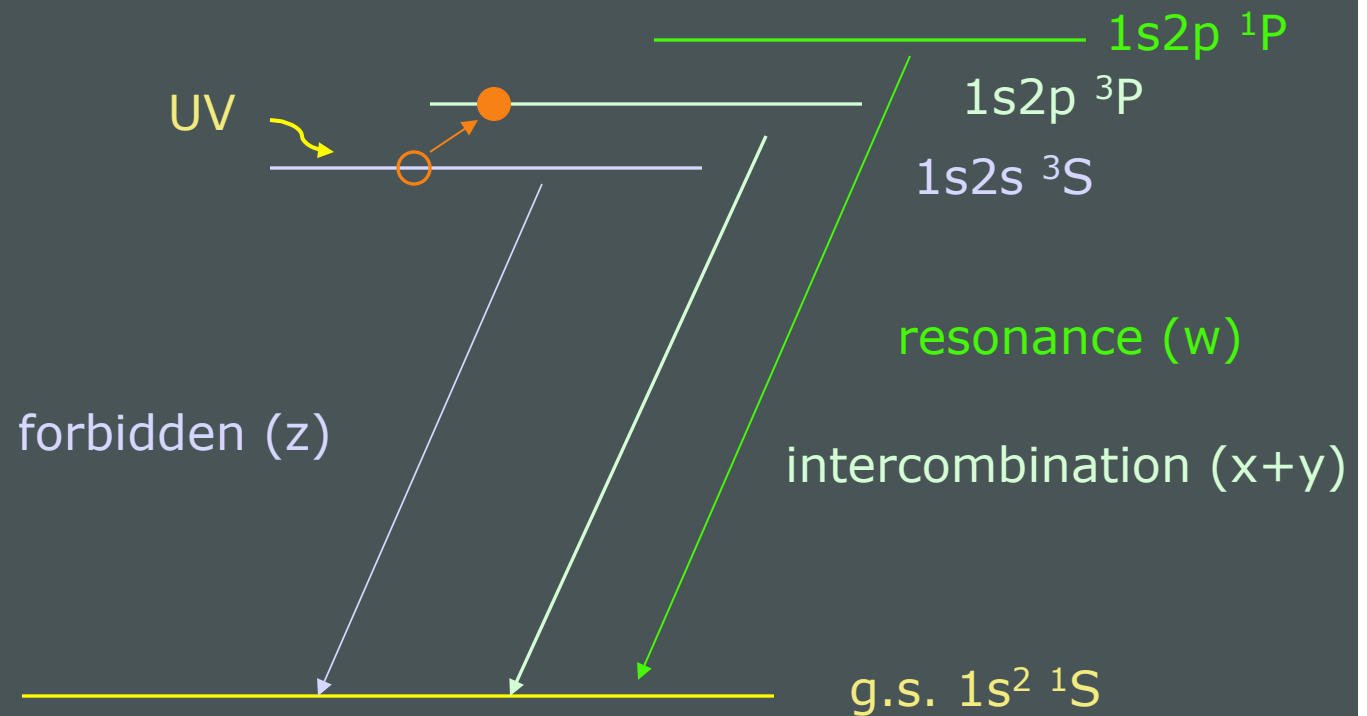
# Helium-like ions (e.g. $O^{+6}$ , $Ne^{+8}$ , $Mg^{+10}$ , $Si^{+12}$ , $S^{+14}$ ) – schematic energy level diagram



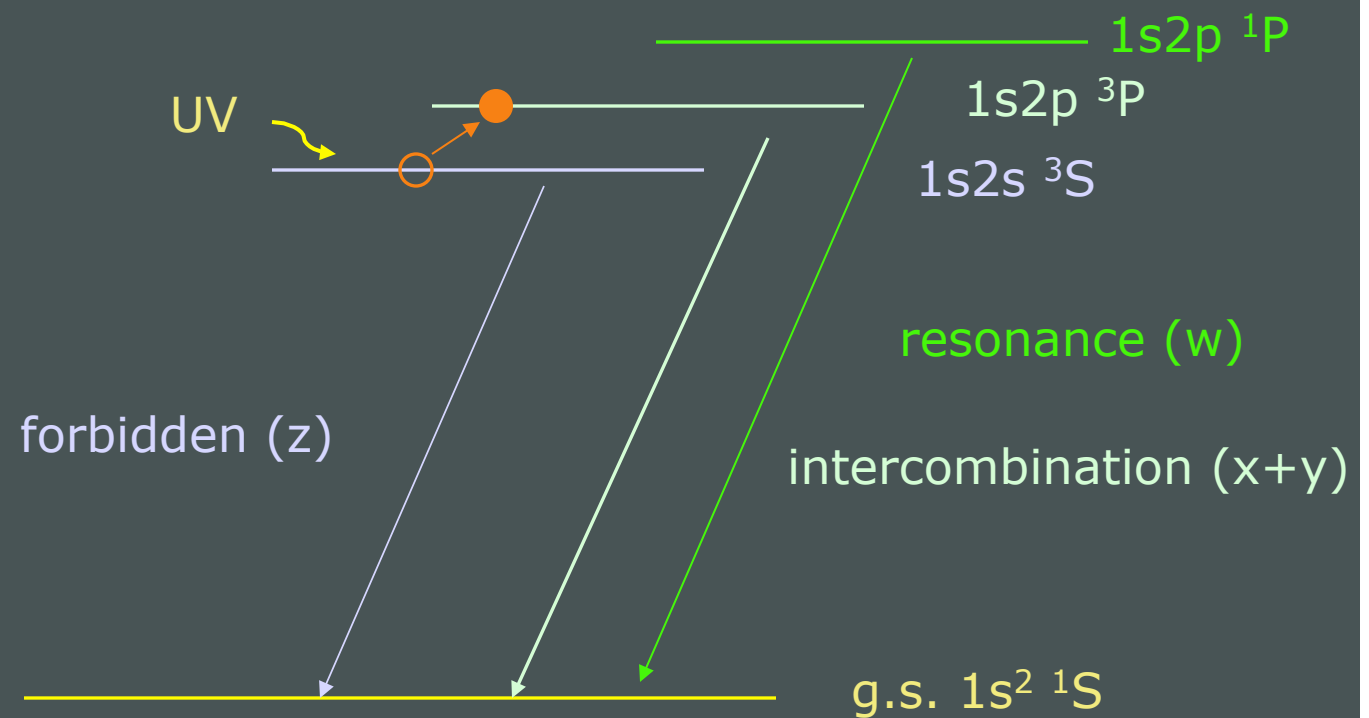
# Ultraviolet light from the star's photosphere drives photoexcitation out of the $^3S$ level



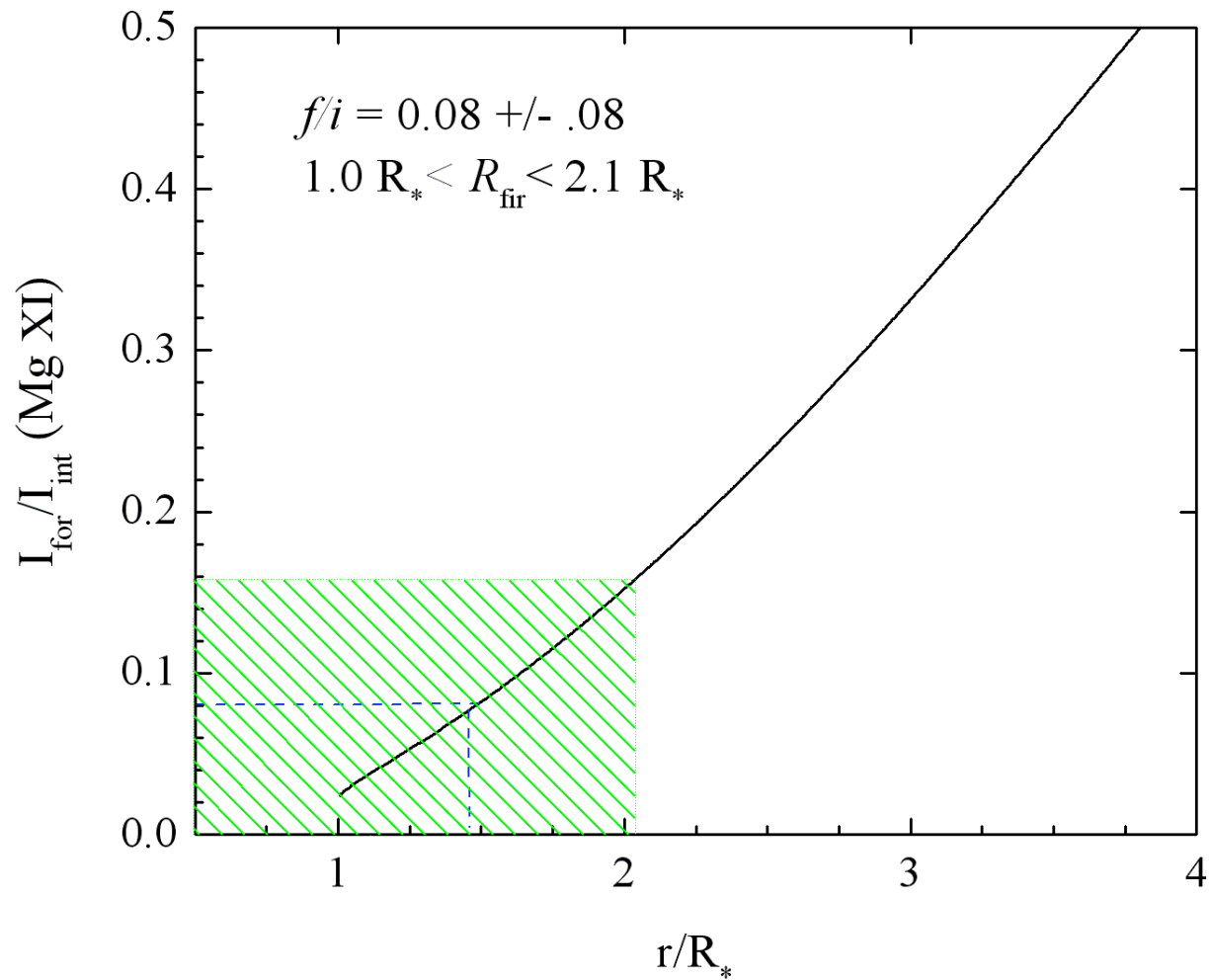
The f/i ratio is thus a diagnostic of the local UV mean intensity...

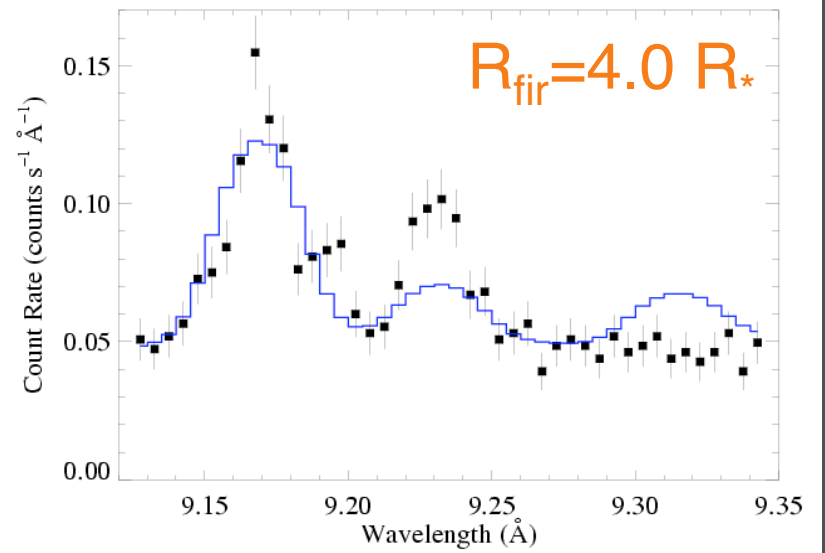
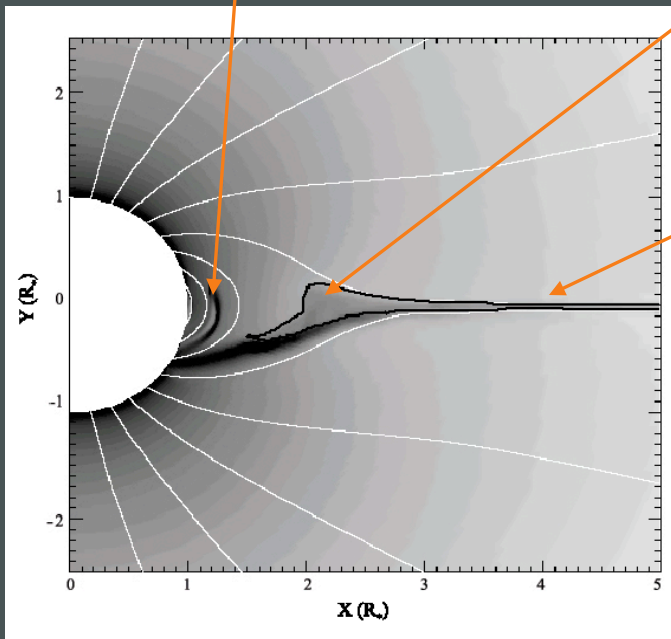
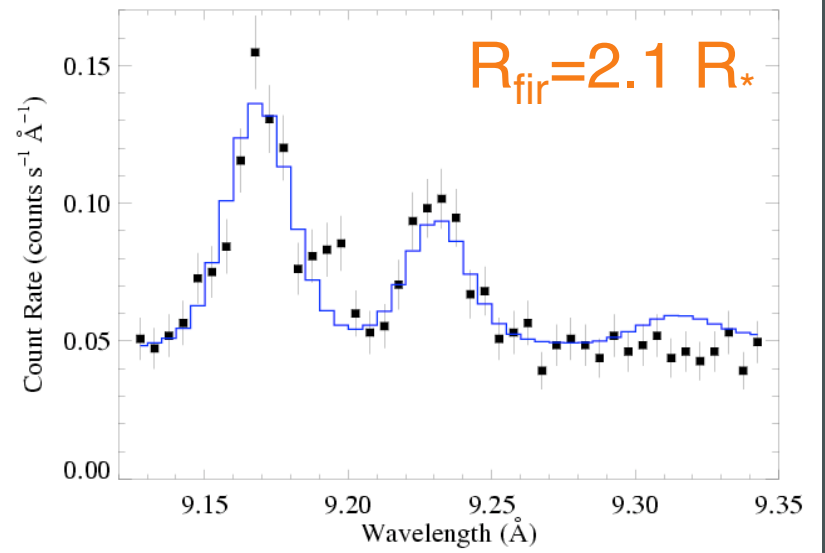
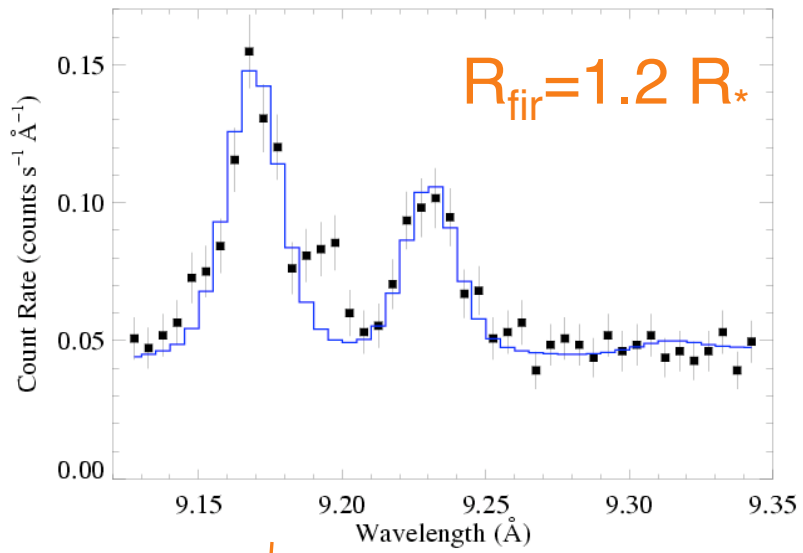


...and thus the distance of the x-ray emitting plasma from the photosphere



# Model of f/i ratio dependence on dilution factor (radius)





He-like f/i ratios and the x-ray light curve both indicate that the hot plasma is somewhat closer to the photosphere than the MHD models predict.

# Conclusions

Normal massive stars have x-ray line profiles consistent with the predictions of the wind instability model.

Photoelectric absorption's effect on the profile shapes can be used as a mass-loss rate diagnostic: *mass-loss rates are lower than previously thought.*

Magnetic massive stars have harder spectra with narrower lines and rotationally modulated variability, in general agreement with MHD simulations.

Line ratio diagnostics are useful for localizing the hot plasma, and indicate that the *MHD simulations predict a location that is too far from the photosphere.*